

Isolated halos as cosmological probes for Dark Energy

Paul de Fromont

supervised by *Jean-Michel Alimi*

July 2, 2014

- 1 The physical context
- 2 Cosmological structures as probes for dark energy
- 3 Isolated halos

Sommaire

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Introduction

Dark energy problem

Today's Universe composition

- Baryonic matter : $\sim 5\%$
- Dark Matter (CDM) : $\sim 25\%$
- Dark Energy : $\sim 70\% \Rightarrow$ accelerated expansion of the Universe

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Dark energy nature ?

- equation of state : $p_{DE} = \omega \rho_{DE}$
- acceleration $\Rightarrow \omega < -1/3$
- in GR, $\omega \in [-1, -1/3]$
- 3 cases :
 - ① $-1 < \omega \leq -1/3$, *quintessence*
 - ② $\omega = -1$: cosmological constant $\rho_{DE} = cste$
 - ③ $\omega < -1$: extension of GR, *modified gravity*

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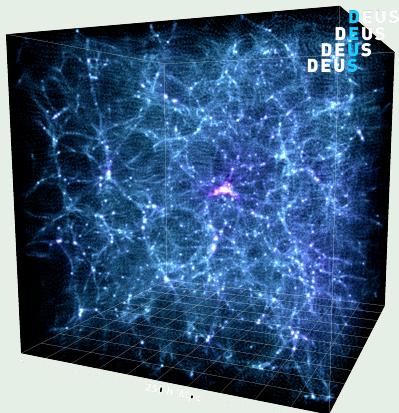
cosmic structures as probes for dark energy

DEUS simulation

N-body simulations

- Λ CDM with $\omega = -1$
- R PCDM quintessence model with $\omega \sim -0.87$
- S UCDM idem with $\omega \sim -0.94$

box of 250 Mpc/h



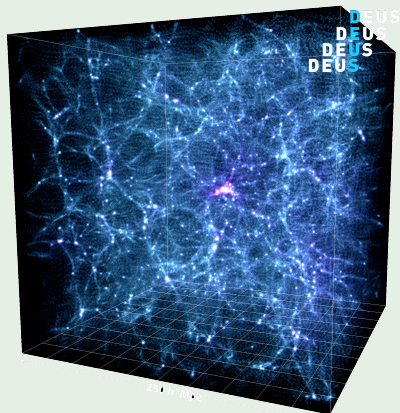
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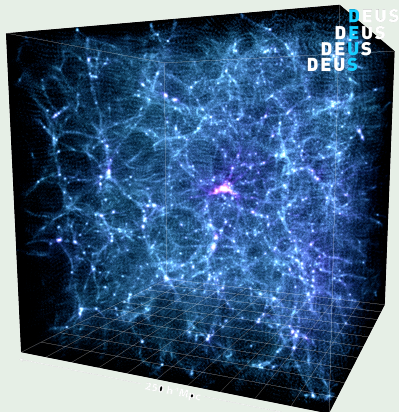
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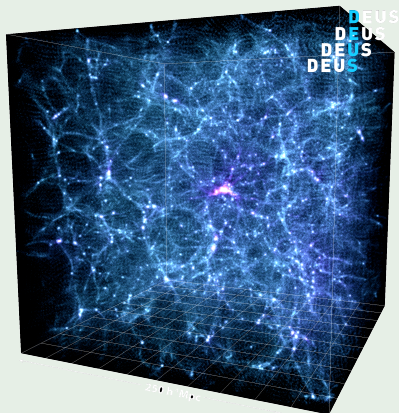
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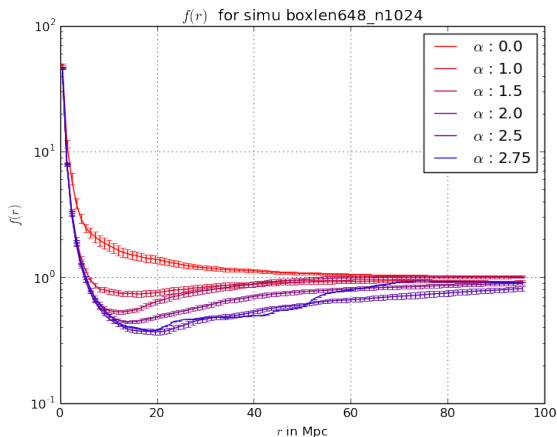
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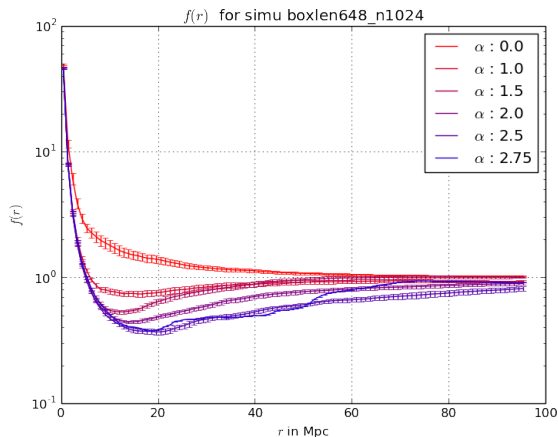


mass profile

$$f(r) = \frac{M(r)}{4/3\pi\bar{\rho}r^3}$$

for halos of mass
 $M_h \sim 5.1 \cdot 10^{12} M_o$
and $d_i \geq \alpha(R_i + R_0)$

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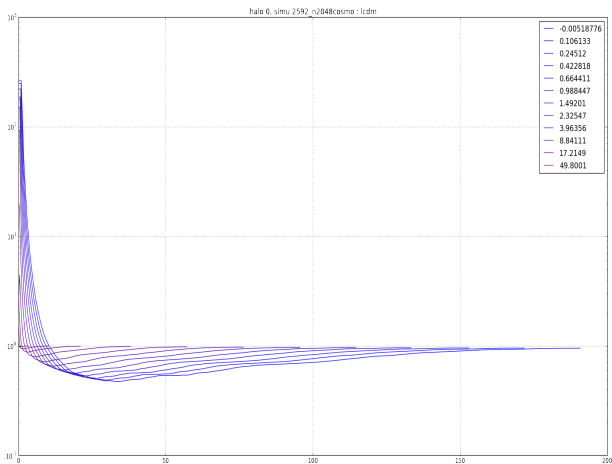


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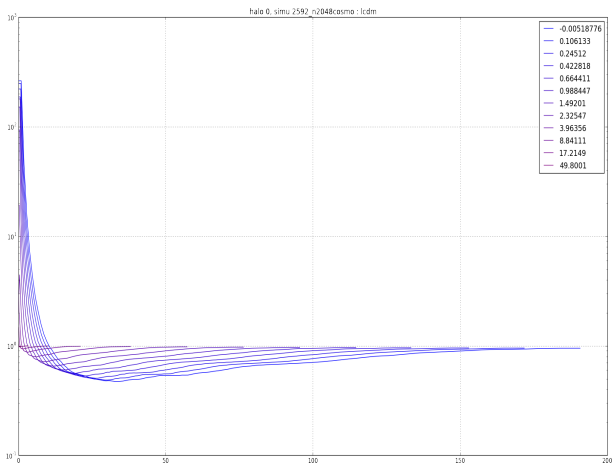
profile evolution



shape evolution

evolution of an initial profile

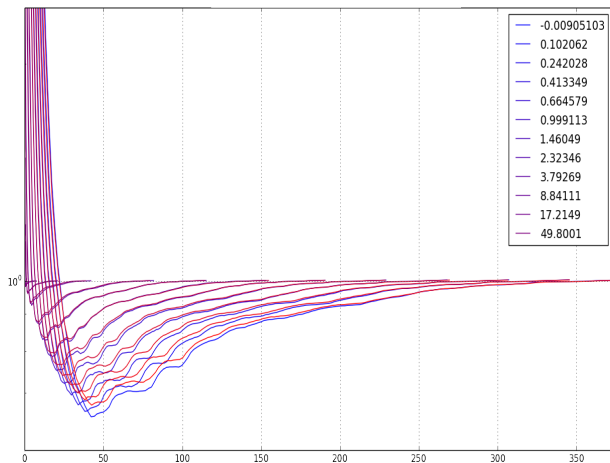
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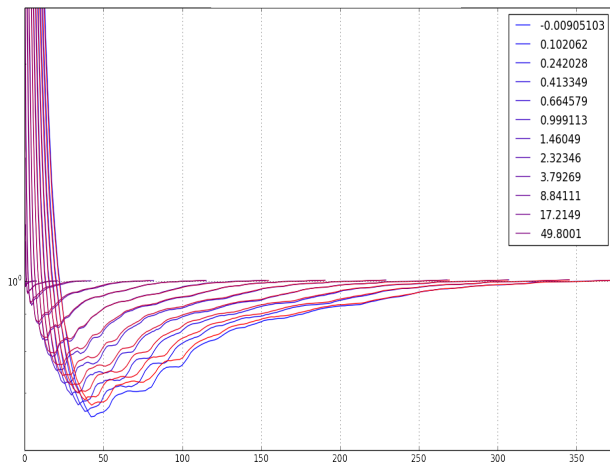
difference between cosmologies



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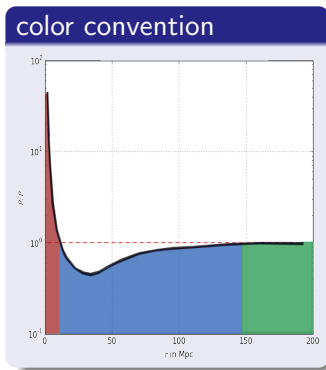
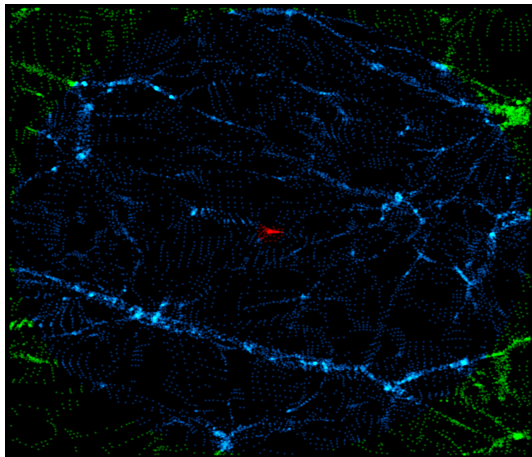
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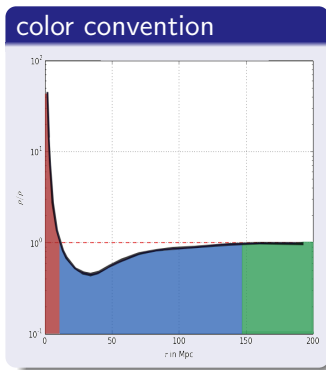
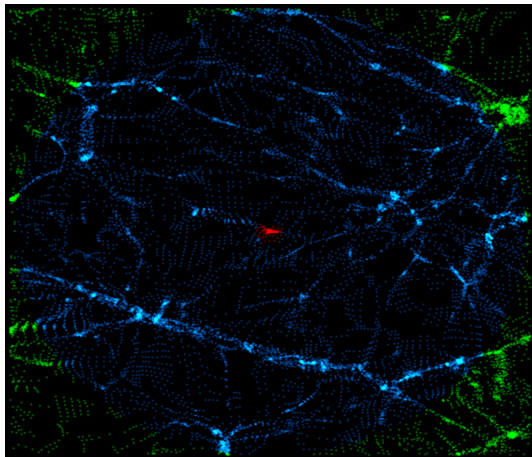


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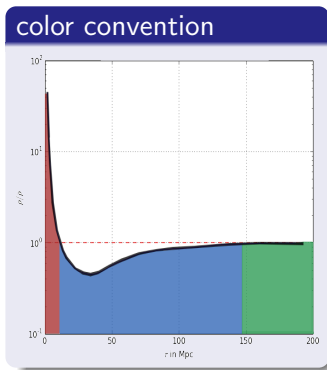
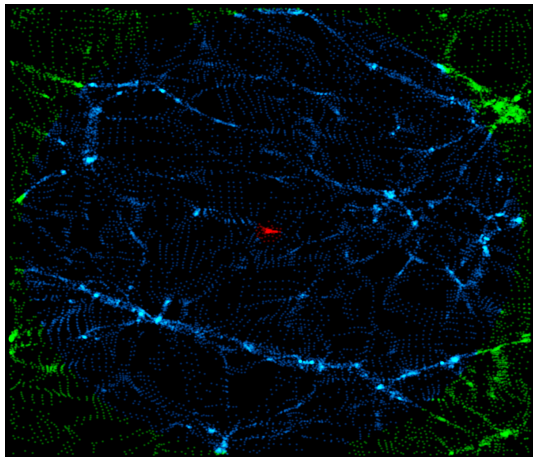
evolution of an isolated halo



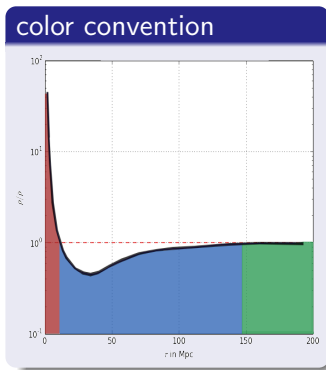
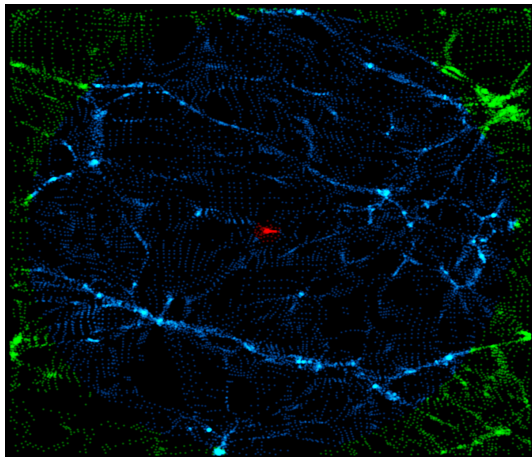
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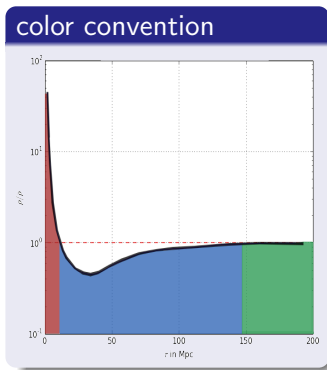
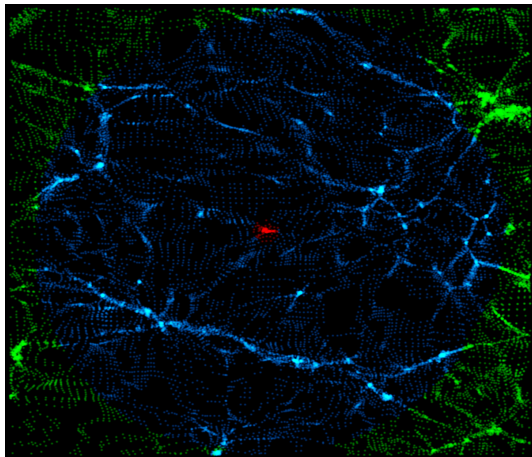
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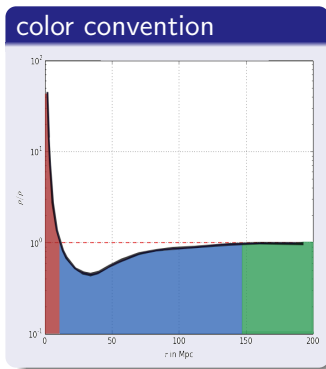
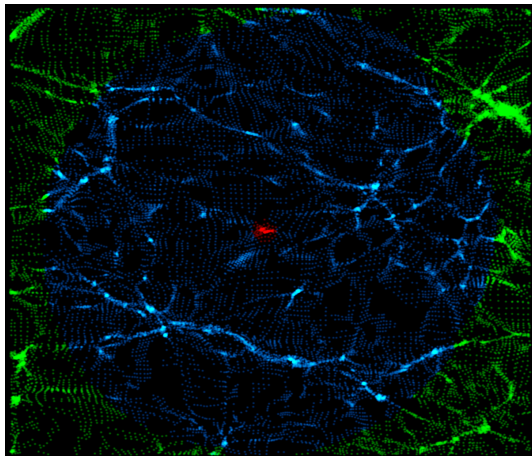
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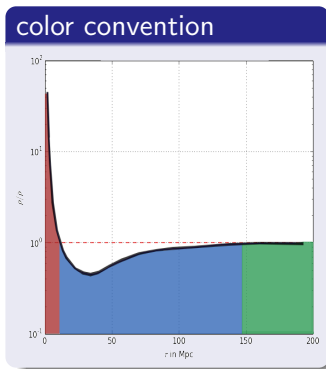
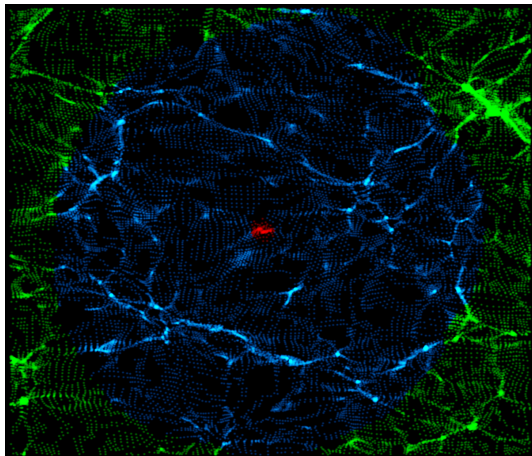
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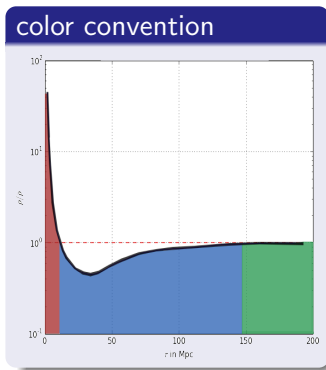
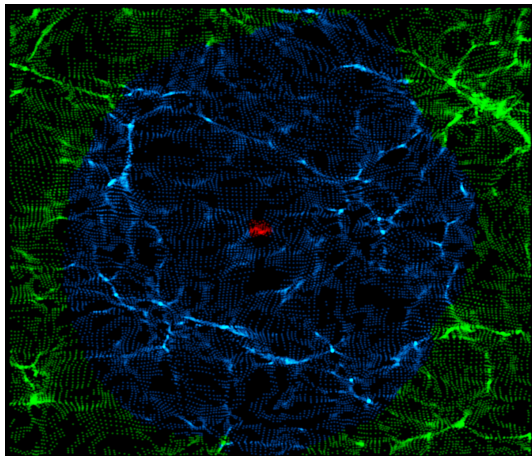
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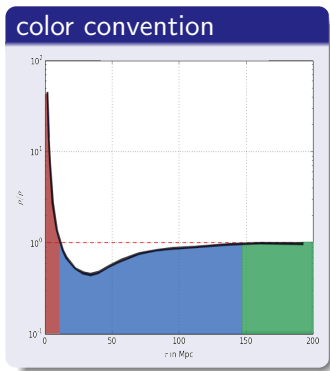
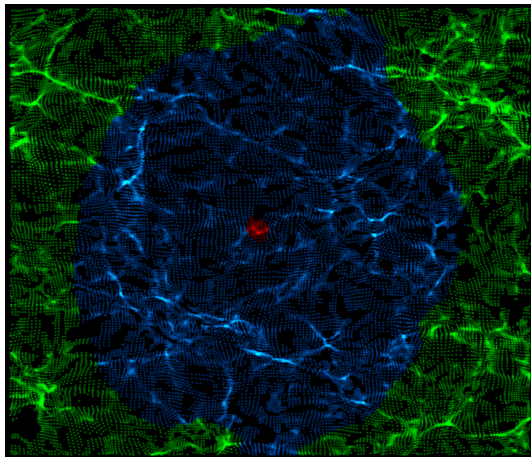
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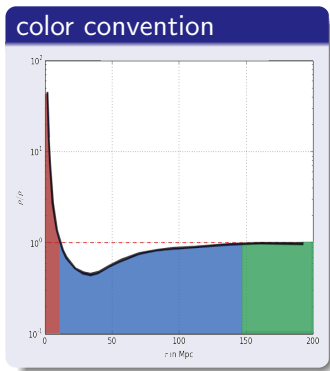
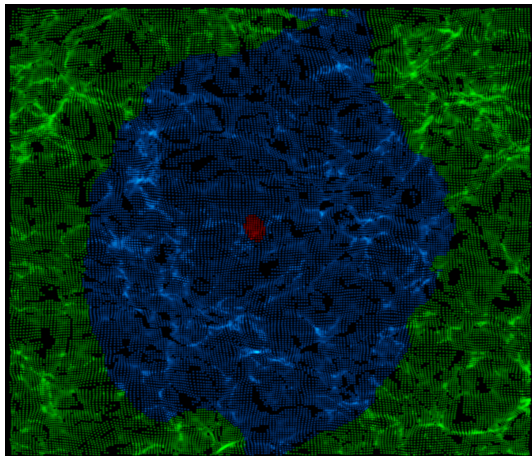
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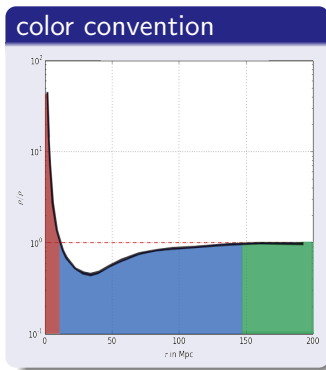
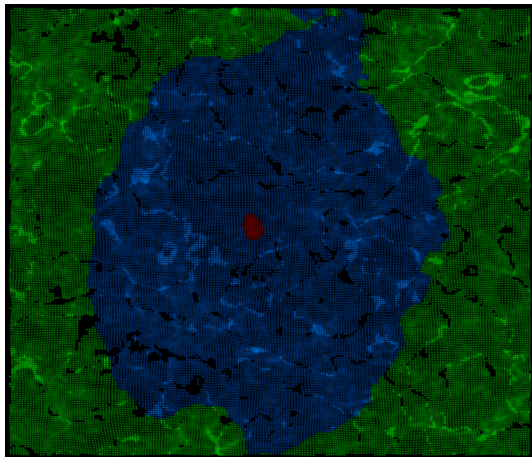
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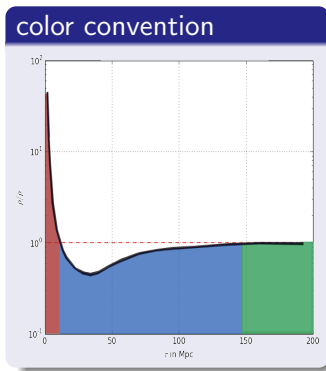
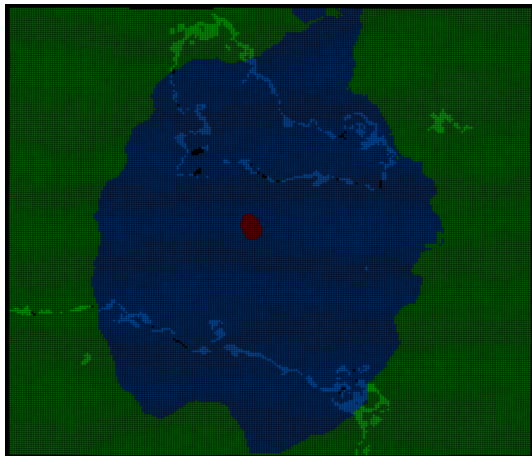
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dynamical evolution

dynamical properties

- no *shell-crossing* !
- "self similar" evolution for $f_0(r) \rightarrow f(r, t)$
- r_1 such as $f(r_1, t) = 1 \Rightarrow r_1(t) \propto a(t)$
- $f_{min}(t)$ such as $\partial_r f(r, t) = 0$

$$f_{min}(t) = \frac{f_{min}(0)}{x(t, f_{min}(0))^3}$$

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dynamical evolution

Evolution equations

introducing $x(t, r_0) = r(t, r_0)/a(t)$

$$\frac{\partial^2 x}{\partial \tau^2} + \frac{1}{\sqrt{2\Omega_m(\tau)}} \frac{\partial x}{\partial \tau} = x - \frac{f_0(r_0)}{x^2}$$

with

$$\tau = \frac{\sqrt{2}}{3\omega} \sinh^{-1} \left[\sqrt{\frac{\Omega_m^0}{1 - \Omega_m^0}} \left(\frac{a(t)}{a(0)} \right)^{3\omega/2} \right]$$

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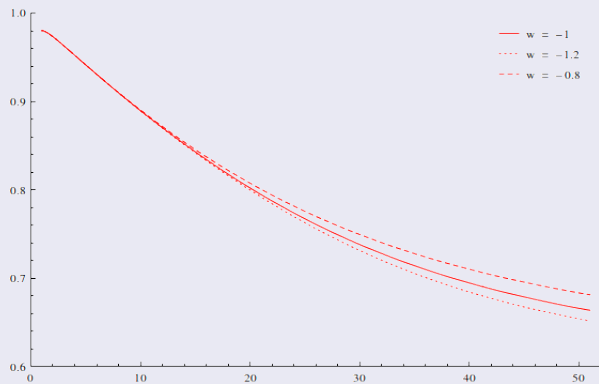
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cosmological dependence of $f_{min}(a)$



Isolated halos

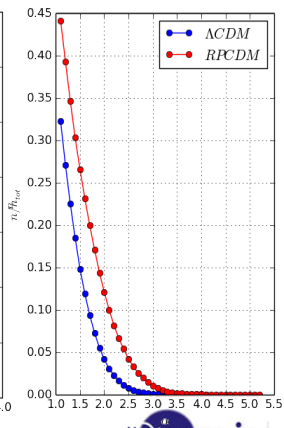
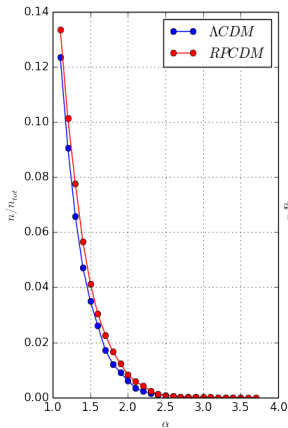
number of isolated halos

for halos of mass :

$M_h \sim 5.2 \cdot 10^{12} M_o$ (left)

$M_h \sim 3.2 \cdot 10^{15} M_o$

(right)



Isolated halos

why useful ?

- known dynamics
- access to $f_0(r)$ (*Gaussian Random Field* and $P(k)$)
- gives $f_{min}(t)$, $r_{min}(t)$...
- probability to be isolated $p(\alpha, M_h)$
- depends on the cosmology !

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Conclusion

Isolated halos

shape and population leads to the cosmology !

work still in progress

finish the job ... include *modified gravity*

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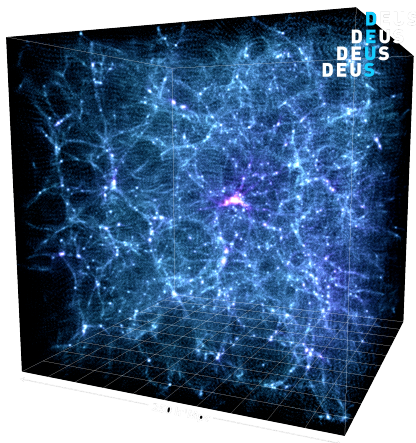
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Thank you for your attention !



Isolated halos

what is an isolated halo ?

No neighbour around ...

$$\forall i, \|\vec{x}_0 - \vec{x}_i\| \geq \alpha(R_i + R_0)$$

- $R_k = \left(\frac{3M_k}{4\pi\rho}\right)^{1/3}$
- $\alpha = 0$ selects all halos

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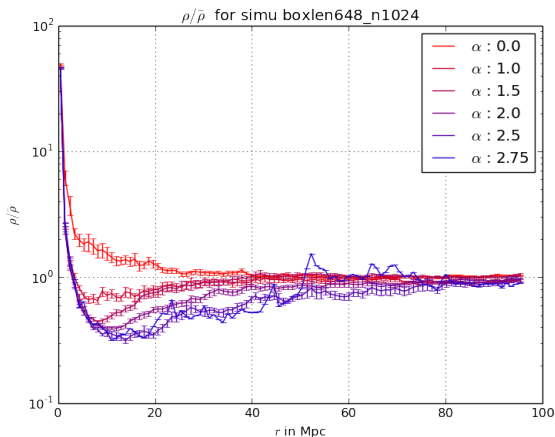
what is an isolated halo ?

No neighbour around ...

$$\forall i, \|\vec{x}_0 - \vec{x}_i\| \geq \alpha(R_i + R_0)$$

- $R_k = \left(\frac{3M_k}{4\pi\rho}\right)^{1/3}$
- $\alpha = 0$ selects all halos

density and mass profile



density profile

$$\delta(r) = \frac{\rho(r)}{\bar{\rho}}$$

for halos of mass
 $M_h \sim 5.1 \cdot 10^{12} M_\odot$