

Gravitational wave modelling and dipolar tidal effects in scalar-tensor theories

Laura BERNARD

Journée LISA à l'Observatoire

vendredi 6 décembre 2021

Going beyond GR

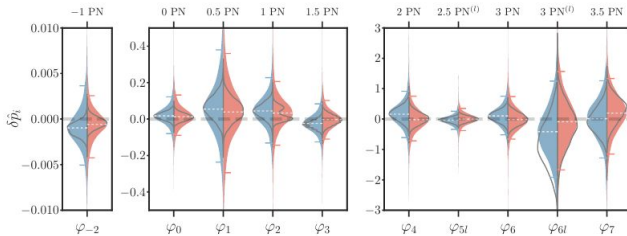
- Why?
 - ▷ dark sectors
 - ▷ quantum gravity
 - ▷ for the beauty

Going beyond GR

- Why?
 - ▷ dark sectors
 - ▷ quantum gravity
 - ▷ for the beauty
- How?
 - ▷ higher dimensions
 - ▷ new fields
 - ▷ etc.

Testing gravity

- ▶ Parametrized vs specific theories tests



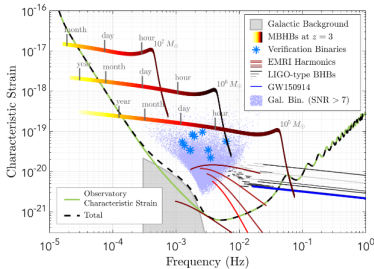
LIGO-Virgo collaboration, 2019

Testing gravity

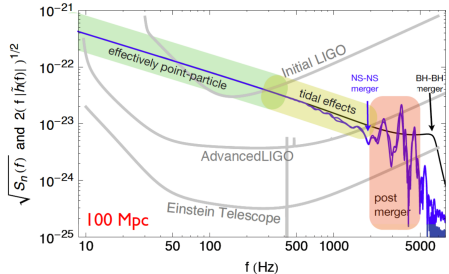
- ▶ Parametrized vs specific theories tests
- ▶ Challenges for modelisation of strong-field effects beyond GR, specially for analytical models
 - *tidal effects, scalarisation, boson clouds, etc.*
 - *what method : EFT, amplitudes, classical PN?*
- ▶ Degeneracies with other effects, ex : *tidal vs eos for NSs*
 - *I-Love test : theory agnostic and EoS insensitive*

Testing gravity

- ▶ Do we really have a chance to be surprised?
 - with LIGO-Virgo, LISA, 3rd generation detectors?
 - using multimessenger astronomy (EHT, NICER)?

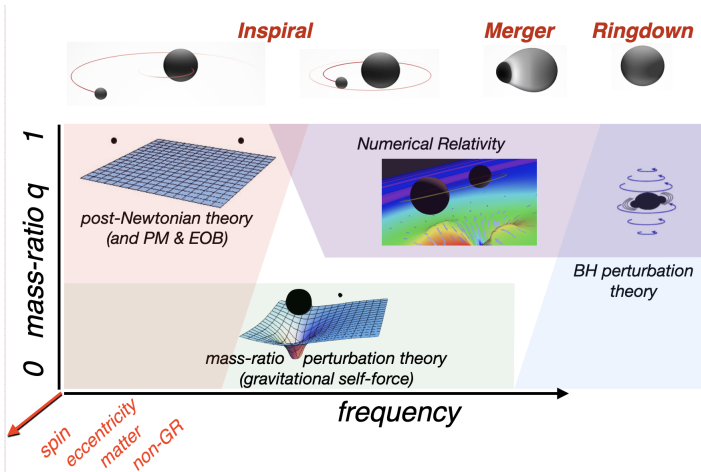


LISA Astro2020 white paper



ET science case

Gravitational wave modelling

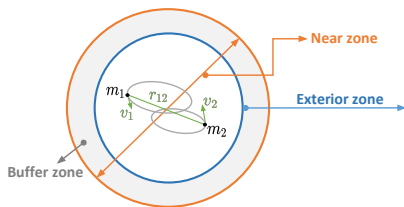


Credits : H. Pfeiffer

Post-Newtonian formalism

Post-Newtonian source

Isolated, compact, slowly moving and weakly stressed source



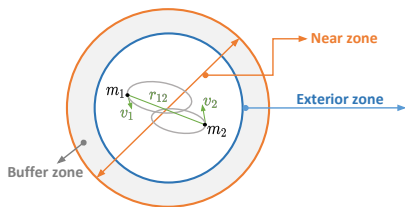
$$\epsilon \equiv \frac{v_{12}^2}{c^2} \sim \frac{Gm}{r_{12}c^2} \ll 1$$

Near zone post-Newtonian expansion, $n\text{PN} = \mathcal{O}\left(\frac{1}{c^{2n}}\right)$

Post-Newtonian formalism

Post-Newtonian source

Isolated, compact, slowly moving and weakly stressed source



$$\epsilon \equiv \frac{v_{12}^2}{c^2} \sim \frac{Gm}{r_{12}c^2} \ll 1$$

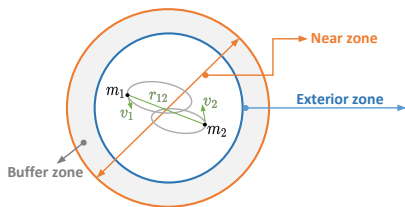
Near zone post-Newtonian expansion, $n\text{PN} = \mathcal{O}\left(\frac{1}{c^{2n}}\right)$

Exterior zone multipolar expansion in power of $\frac{r_{12}}{R}$

Post-Newtonian formalism

Post-Newtonian source

Isolated, compact, slowly moving and weakly stressed source



$$\epsilon \equiv \frac{v_{12}^2}{c^2} \sim \frac{Gm}{r_{12}c^2} \ll 1$$

Near zone post-Newtonian expansion, $n\text{PN} = \mathcal{O}\left(\frac{1}{c^{2n}}\right)$

Exterior zone multipolar expansion in power of $\frac{r_{12}}{R}$

Matching radiative moments \longleftarrow source moments \longrightarrow source
exp. in $1/R$ matching

Scalar-tensor theories

$$S_{\text{ST}} = \frac{c^3}{16\pi G} \int d^4x \sqrt{-g} \left[\phi R - \frac{\omega(\phi)}{\phi} g^{\alpha\beta} \partial_\alpha \partial_\beta \phi \right] + S_m(\mathbf{m}, g_{\alpha\beta})$$

- ▶ well-posed, passes solar system tests
- ▶ no hair theorem
- ▶ neutron stars : scalarization

Coupling to matter

Violation of the Strong Equivalence Principle

- Incorporate the internal structure of compact, self-gravitating bodies
- Eardley's approach : masses depend on the scalar field $m_A(\phi)$

$$S_m = -c \sum_A \int d\tau_A m_A(\phi)$$

Coupling to matter

Violation of the Strong Equivalence Principle

- Incorporate the internal structure of compact, self-gravitating bodies
- Eardley's approach : masses depend on the scalar field $m_A(\phi)$

$$S_m = -c \sum_A \int d\tau_A m_A(\phi)$$

- ▷ Sensitivities : $s_A = \left. \frac{d \ln m_A(\phi)}{d \ln \phi} \right|_0$
- Neutron stars : $s_A \sim 0.2$ (depends on the equation of states)
 - Black holes : $s_A = 0.5$ (compactness M/R)

Equations of motion

$$\frac{d\mathbf{v}_1}{dt} = \underbrace{-\frac{G_{\text{eff}} m_2}{r_{12}^2} \mathbf{n}_{12} + \frac{\mathbf{A}_{1\text{PN}}}{c^2}}_{\text{conservative terms}} + \underbrace{\frac{\mathbf{A}_{2\text{PN}}}{c^4}}_{\text{cons.}}$$

Equations of motion

$$\frac{d\mathbf{v}_1}{dt} = \underbrace{-\frac{G_{\text{eff}} m_2}{r_{12}^2} \mathbf{n}_{12} + \frac{\mathbf{A}_{1\text{PN}}}{c^2}}_{\text{conservative terms}} + \underbrace{\frac{\mathbf{A}_{1.5\text{PN}}}{c^3}}_{\text{rad. reac.}} + \underbrace{\frac{\mathbf{A}_{2\text{PN}}}{c^4}}_{\text{cons.}} + \underbrace{\frac{\mathbf{A}_{2.5\text{PN}}}{c^5}}_{\text{rad. reac.}}$$

Differences w.r.t. GR

- Dissipative effects start at 1.5PN

Equations of motion

$$\begin{aligned}
 \frac{d\mathbf{v}_1}{dt} = & \underbrace{-\frac{G_{\text{eff}} m_2}{r_{12}^2} \mathbf{n}_{12} + \frac{\mathbf{A}_{1\text{PN}}}{c^2}}_{\text{conservative terms}} + \underbrace{\frac{\mathbf{A}_{1.5\text{PN}}}{c^3}}_{\text{rad. reac.}} + \underbrace{\frac{\mathbf{A}_{2\text{PN}}}{c^4}}_{\text{cons.}} + \underbrace{\frac{\mathbf{A}_{2.5\text{PN}}}{c^5}}_{\text{rad. reac.}} \\
 & + \underbrace{\frac{\mathbf{A}_{3\text{PN}}^{\text{inst}}}{c^6}}_{\text{cons, local}} + \underbrace{\frac{\mathbf{A}_{3\text{PN}}^{\text{tail}}}{c^6}}_{\text{cons, nonloc.}}
 \end{aligned}$$

Differences w.r.t. GR

- Dissipative effects start at 1.5PN
- A conservative scalar tail term at 3PN : $\mathbf{A}_{3\text{PN}}^{\text{tail}} \propto \int_{-\infty}^{+\infty} \frac{dt'}{|t-t'|} I_i^{(4)}(t')$

Equations of motion

$$\begin{aligned}
 \frac{d\mathbf{v}_1}{dt} = & \underbrace{-\frac{G_{\text{eff}} m_2}{r_{12}^2} \mathbf{n}_{12} + \frac{\mathbf{A}_{1\text{PN}}}{c^2}}_{\text{conservative terms}} + \underbrace{\frac{\mathbf{A}_{1.5\text{PN}}}{c^3}}_{\text{rad. reac.}} + \underbrace{\frac{\mathbf{A}_{2\text{PN}}}{c^4}}_{\text{cons.}} + \underbrace{\frac{\mathbf{A}_{2.5\text{PN}}}{c^5}}_{\text{rad. reac.}} \\
 & + \underbrace{\frac{\mathbf{A}_{3\text{PN}}^{\text{inst}}}{c^6}}_{\text{cons, local}} + \underbrace{\frac{\mathbf{A}_{3\text{PN}}^{\text{tail}}}{c^6}}_{\text{cons, nonloc.}} + \underbrace{\frac{\mathbf{A}_{3\text{PN}}^{\text{tidal}}}{c^6}}_{\text{cons, local}}
 \end{aligned}$$

Differences w.r.t. GR

- Dissipative effects start at 1.5PN
- A conservative scalar tail term at 3PN : $\mathbf{A}_{3\text{PN}}^{\text{tail}} \propto \int_{-\infty}^{+\infty} \frac{dt'}{|t-t'|} I_i^{(4)}(t')$
- Tidal effects start at 3PN

Scalar and gravitational flux

$$\mathcal{F} = \frac{32c^5\nu^2x^5}{5G_{\text{eff}}} \left[1 + \frac{\mathcal{F}_{1\text{PN}}^{\text{grav}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{grav}}}{c^3} \right]$$

Scalar and gravitational flux

$$\mathcal{F} = \frac{32c^5\nu^2x^5}{5G_{\text{eff}}} \left[1 + \frac{\mathcal{F}_{1\text{PN}}^{\text{grav}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{grav}}}{c^3} \right] + \frac{4c^5\nu^2x^5}{3G_{\text{eff}}} \zeta S_-^2 \left[x^{-1} \right]$$

Differences w.r.t. GR

- Scalar flux starts at -1PN

Scalar and gravitational flux

$$\mathcal{F} = \frac{32c^5\nu^2x^5}{5G_{\text{eff}}} \left[1 + \frac{\mathcal{F}_{1\text{PN}}^{\text{grav}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{grav}}}{c^3} \right] \\ + \frac{4c^5\nu^2x^5}{3G_{\text{eff}}} \zeta S_-^2 \left[x^{-1} + \frac{\mathcal{F}_{0\text{PN}}^{\text{scal}}}{c^0} + \frac{\mathcal{F}_{0.5\text{PN}}^{\text{scal}}}{c^1} + \frac{\mathcal{F}_{1\text{PN}}^{\text{scal}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{scal}}}{c^3} \right]$$

Differences w.r.t. GR

- Scalar flux starts at -1PN
- A scalar tail term at 0.5PN $\delta U_{ij} \propto \int \frac{dt'}{|t-t'|} I_s^{(2)}(t') I_{sj}^{(2)}(t')$
- A scalar memory term at 1.5PN $\delta U_i^s \propto M \int d\tau \ln\left(\frac{\tau}{\tau_0}\right) I_{si}^{(3)}(t-\tau)$

Scalar and gravitational flux

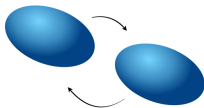
$$\mathcal{F} = \frac{32c^5\nu^2x^5}{5G_{\text{eff}}} \left[1 + \frac{\mathcal{F}_{1\text{PN}}^{\text{grav}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{grav}}}{c^3} \right] + \frac{4c^5\nu^2x^4}{3G_{\text{eff}}} \zeta S_-^2 \cdot \frac{\mathcal{F}_{2\text{PN}}^{\text{scal, tidal}}}{c^4}$$

$$+ \frac{4c^5\nu^2x^5}{3G_{\text{eff}}} \zeta S_-^2 \left[x^{-1} + \frac{\mathcal{F}_{0\text{PN}}^{\text{scal}}}{c^0} + \frac{\mathcal{F}_{0.5\text{PN}}^{\text{scal}}}{c^1} + \frac{\mathcal{F}_{1\text{PN}}^{\text{scal}}}{c^2} + \frac{\mathcal{F}_{1.5\text{PN}}^{\text{scal}}}{c^3} \right]$$

Differences w.r.t. GR

- Scalar flux starts at -1PN
- A scalar tail term at 0.5PN $\delta U_{ij} \propto \int \frac{dt'}{|t-t'|} I_s^{(2)}(t') I_{sj}^{(2)}(t')$
- A scalar memory term at 1.5PN $\delta U_i^s \propto M \int d\tau \ln\left(\frac{\tau}{\tau_0}\right) I_{si}^{(3)}(t-\tau)$
- A scalar tidal contribution at 2PN

Tidal effects - Newtonian theory

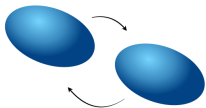


$$\mathcal{E}_{ij} \sim \partial_i \partial_j U$$

Response to an external tidal field \mathcal{E}_{ij}

$$U = \frac{M}{R} - \frac{1}{2} \mathcal{E}_{ij} x^i x^j + \frac{3}{2} \frac{Q_{ij} x^i x^j}{r^5}$$

Tidal effects - Newtonian theory



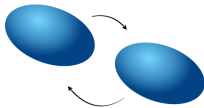
$$\mathcal{E}_{ij} \sim \partial_i \partial_j U$$

Response to an external tidal field \mathcal{E}_{ij}

$$U = \frac{M}{R} - \frac{1}{2} \mathcal{E}_{ij} x^i x^j + \frac{3}{2} \frac{Q_{ij} x^i x^j}{r^5}$$

- ▶ Adiabatical approximation : $Q_{ij} = -\frac{2}{3} k_2 R^5 \mathcal{E}_{ij}$

Tidal effects - Newtonian theory



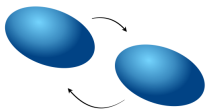
$$\mathcal{E}_{ij} \sim \partial_i \partial_j U$$

Response to an external tidal field \mathcal{E}_{ij}

$$U = \frac{M}{R} - \frac{1}{2} \left(1 + 2k_2 \frac{R^5}{r^5} \right) \mathcal{E}_{ij} x^i x^j$$

- ▶ Adiabatical approximation : $Q_{ij} = -\frac{2}{3} k_2 R^5 \mathcal{E}_{ij}$

Tidal effects - Newtonian theory

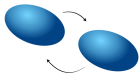


$$\mathcal{E}_{ij} \sim \partial_i \partial_j U$$

Response to an external tidal field \mathcal{E}_{ij}

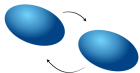
$$U = \frac{M}{R} - \frac{1}{2} \left(1 + 2k_2 \frac{R^5}{r^5} \right) \mathcal{E}_{ij} x^i x^j$$

- ▶ Adiabatical approximation : $Q_{ij} = -\frac{2}{3} k_2 R^5 \mathcal{E}_{ij}$
- ▶ We can generalize to multipolar field \mathcal{E}_L and tidal Love number k_L



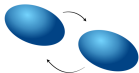
Tidal effects - General relativity

- Electric-type multipole moments : $\mathcal{E}_L \propto \nabla_{L-2} C_{0a_1 0a_2}$
- Magnetic-type multipole moments : $\mathcal{B}_L \propto \epsilon_{a_1 bc} \nabla_{L-2} C_{a_2 0bc}$



Tidal effects - General relativity

- Electric-type multipole moments : $\mathcal{E}_L \propto \nabla_{L-2} C_{0a_1 0a_2}$
- Magnetic-type multipole moments : $\mathcal{B}_L \propto \epsilon_{a_1 bc} \nabla_{L-2} C_{a_2 0bc}$
- ▷ Electric and magnetic tidal Love numbers k_L and j_L



Tidal effects - General relativity

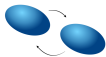
- Electric-type multipole moments : $\mathcal{E}_L \propto \nabla_{L-2} C_{0a_1 0a_2}$
- Magnetic-type multipole moments : $\mathcal{B}_L \propto \epsilon_{a_1 bc} \nabla_{L-2} C_{a_2 0 bc}$
- ▷ Electric and magnetic tidal Love numbers k_L and j_L

In the matter action

$$S_m = S_{\text{pp}} + \sum_A \int d\tau_A [\mu_A \mathcal{E}_{\mu\nu}^A \mathcal{E}_A^{\mu\nu} + \sigma_A \mathcal{B}_{\mu\nu}^A \mathcal{B}_A^{\mu\nu} + \dots]$$

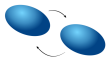
- ▷ In the phase : $\Delta\psi_{\text{fs}} \propto \psi_{\text{N}} \cdot k_2 \cdot \frac{R^5}{r^5}$
- ▷ Formally a 5PN effect but $\propto R^5$

Tidal effects - Scalar-tensor theory



- ▶ Scalar-type multipole moments : $\mathcal{E}_L^{(s)} \sim \nabla_L \varphi$

Tidal effects - Scalar-tensor theory

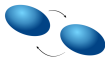


- ▶ Scalar-type multipole moments : $\mathcal{E}_L^{(s)} \sim \nabla_L \varphi$

Response to an external scalar dipolar field

$$U = \frac{M}{R} - \mathcal{E}_i^{(s)} x^i + \frac{Q_i x^i}{r^3}$$

Tidal effects - Scalar-tensor theory



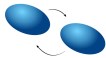
- ▶ Scalar-type multipole moments : $\mathcal{E}_L^{(s)} \sim \nabla_L \varphi$

Response to an external scalar dipolar field

$$U = \frac{M}{R} - \mathcal{E}_i^{(s)} x^i + \frac{Q_i x^i}{r^3}$$

- Adiabatic approximation : $Q_\mu^{(s)} = -\lambda_{(s)} \mathcal{E}_\mu^{(s)}$

Tidal effects - Scalar-tensor theory

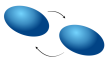


- ▶ Scalar-type multipole moments : $\mathcal{E}_L^{(s)} \sim \nabla_L \varphi$

Response to an external scalar dipolar field

$$U = \frac{M}{R} - \left(1 + k_s \frac{R^3}{r^3} \right) \mathcal{E}_i^{(s)} x^i$$

- Adiabatic approximation : $Q_\mu^{(s)} = -\lambda_{(s)} \mathcal{E}_\mu^{(s)}$
- Scalar-type Love number : $k_s \propto \lambda_s R^3$

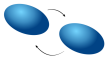


From point particles to extended body

$$m_A(\phi) \longrightarrow m_A[\phi] = m_A(\phi) + N_A(\phi) \nabla_\mu \phi \nabla^\mu \phi$$

In the action

$$S_m = S_{pp} - \frac{1}{2} \sum_A \lambda_A^{(s)} \int d\tau_A (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi)_A + \text{high. orders}$$



From point particles to extended body

$$m_A(\phi) \longrightarrow m_A[\phi] = m_A(\phi) + N_A(\phi) \nabla_\mu \phi \nabla^\mu \phi$$

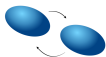
In the action

$$S_m = S_{\text{pp}} - \frac{1}{2} \sum_A \lambda_A^{(s)} \int d\tau_A (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi)_A + \text{high. orders}$$

Consequence on the dynamics

$$\Delta \mathbf{a}_{(fs)} \propto \mathbf{a}_{(N)} \cdot \left[\frac{m_2}{m_1} \bar{\delta}_1 \lambda_1^{(s)} + \frac{m_1}{m_2} \bar{\delta}_2 \lambda_2^{(s)} \right] \frac{1}{r^3}$$

- ▶ formally 3PN order correction with small ST parameters



From point particles to extended body

$$m_A(\phi) \longrightarrow m_A[\phi] = m_A(\phi) + N_A(\phi) \nabla_\mu \phi \nabla^\mu \phi$$

In the action

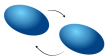
$$S_m = S_{\text{pp}} - \frac{1}{2} \sum_A \lambda_A^{(s)} \int d\tau_A (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi)_A + \text{high. orders}$$

Consequence on the dynamics

$$\Delta \mathbf{a}_{(fs)} \propto \mathbf{a}_{(N)} \cdot \left[\frac{m_2}{m_1} \bar{\delta}_1 k_1^{(s)} + \frac{m_1}{m_2} \bar{\delta}_2 k_2^{(s)} \right] \frac{R^3}{r^3}$$

▷ formally 3PN order correction with small ST parameters

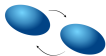
▷ but scales as $\left(\frac{R}{M}\right)^3$



Effect on the gravitational signal

$$\text{Phase evolution : } \frac{d\psi}{dx} = -\frac{(c^2 x)^{3/2}}{\tilde{G}\alpha m} \frac{dE/dx}{\mathcal{F}}$$

$$\frac{\mathcal{F}_{\text{dip}}}{\mathcal{F}_{\text{quad}}} \propto \frac{(s_1 - s_2)^2}{x}$$

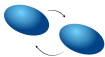


Effect on the gravitational signal

$$\frac{\mathcal{F}_{\text{dip}}}{\mathcal{F}_{\text{quad}}} \propto \frac{(s_1 - s_2)^2}{x}$$

Quadrupolar-driven regime

$$\Delta\psi_{(f_s)} \propto -\frac{1}{32\zeta\eta x^{5/2}} k_s \frac{R^3}{r^3} \implies \text{non detectable}$$



Effect on the gravitational signal

$$\frac{\mathcal{F}_{\text{dip}}}{\mathcal{F}_{\text{quad}}} \propto \frac{(s_1 - s_2)^2}{x}$$

Quadrupolar-driven regime

$$\Delta\psi_{(f_s)} \propto -\frac{1}{32\zeta\eta x^{5/2}} k_s \frac{R^3}{r^3} \implies \text{non detectable}$$

Dipolar-driven regime

$$\Delta\psi_{(f_s)} \propto -\frac{1}{(s_1 - s_2)^2 \eta x^{7/2}} k_s \frac{R^3}{r^3}$$

- ▶ formally 2PN effect in the phase (beyond GR)
- ▶ but similar to the ST 1PN contribution
- ▶ may contribute $\mathcal{O}(1)$ cycles \implies detectable by LISA or 3G

Conclusion

State-of-the-art in ST theories

- Equations of motion at 3PN order
- Total flux and scalar modes at 1.5PN order
- Scalar tidal effect at 3PN order

Conclusion

State-of-the-art in ST theories

- Equations of motion at 3PN order
- Total flux and scalar modes at 1.5PN order
- Scalar tidal effect at 3PN order

To be done

- ▷ Total flux and scalar modes at 2.5PN order
- ▷ Generalise to higher multipolar scalar tides
- ▷ Spin effects \implies other theories?