



Propagating light through the large-scale structure of the Universe

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16/12/2021

Our universe

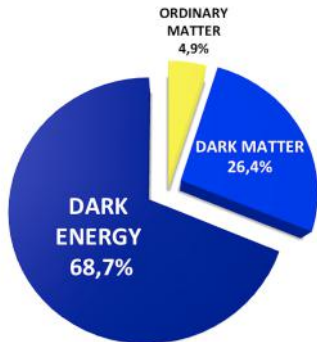
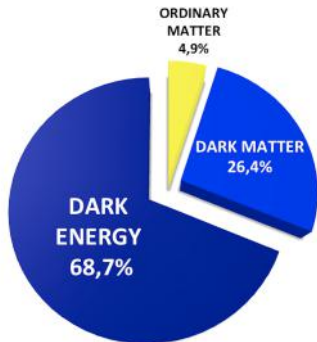


FIGURE – Planck (2018) assuming Λ CDM

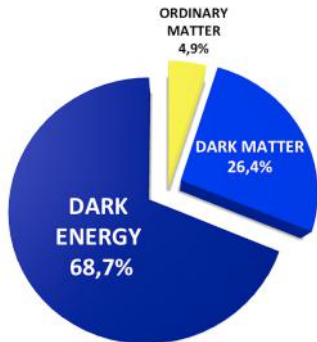
Our universe



- Nature of dark sector is unknown
- Dark sector not directly observable

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Indirect observations :

- Galaxy rotation curves
- CMB
- etc...

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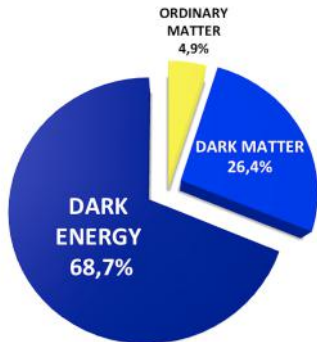


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Future surveys will observe billions of galaxies

Need additional probes or refine existing ones

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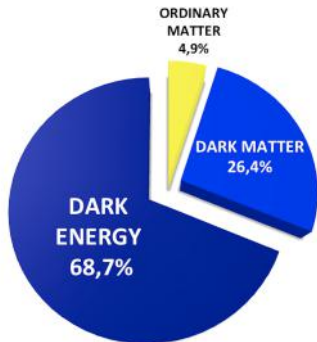


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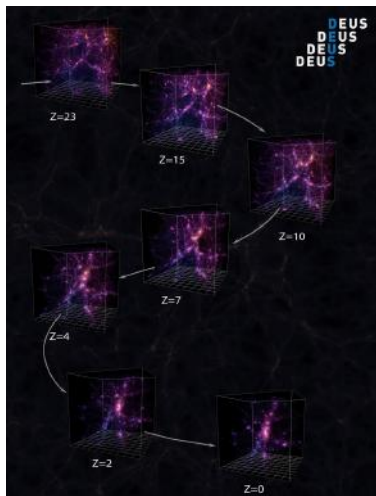
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Question

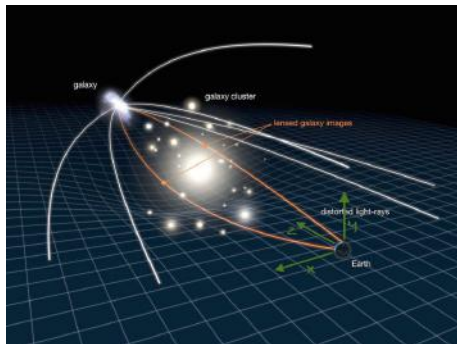
What are these dark matter and dark energy components which dominate the energy content of our universe ?

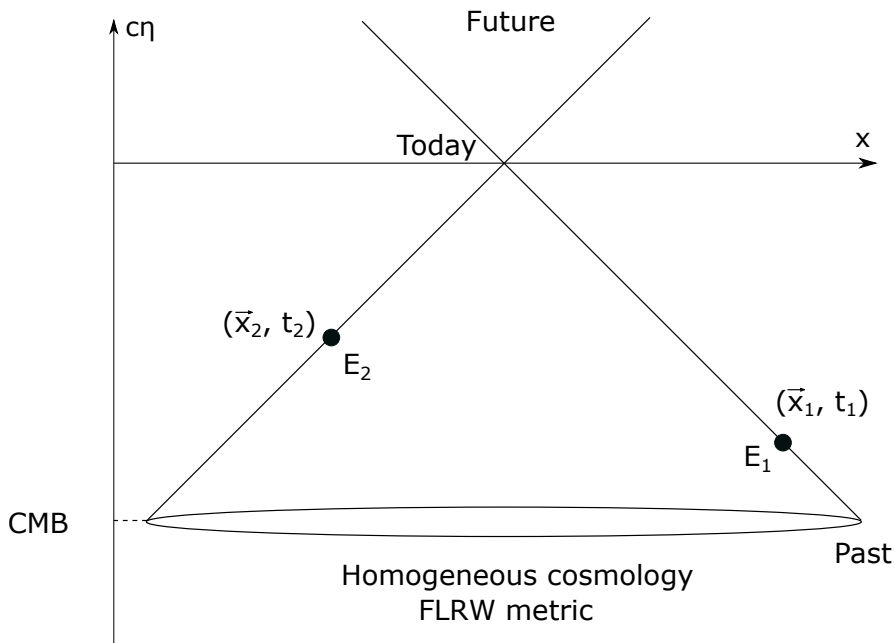
Impact of the dark components on

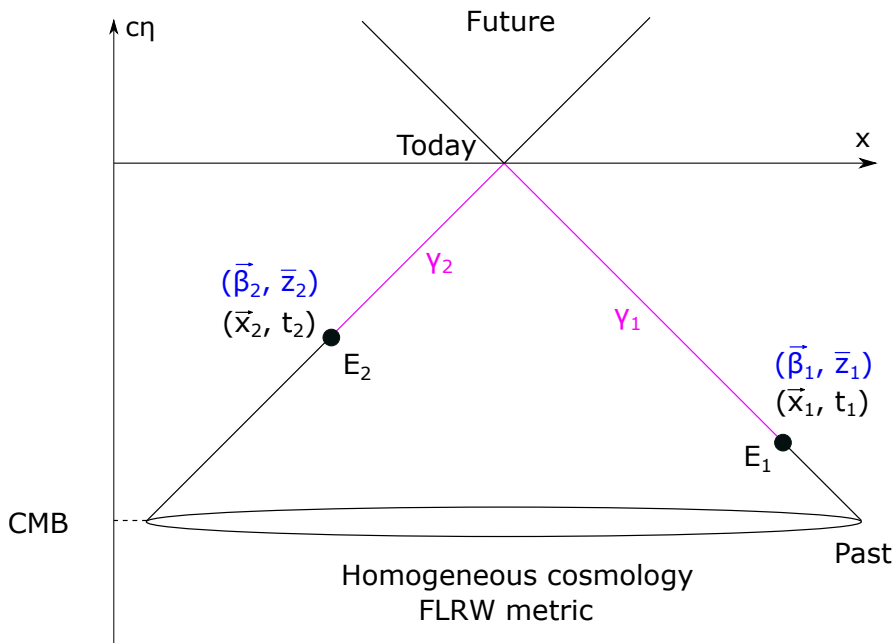
Structure formation

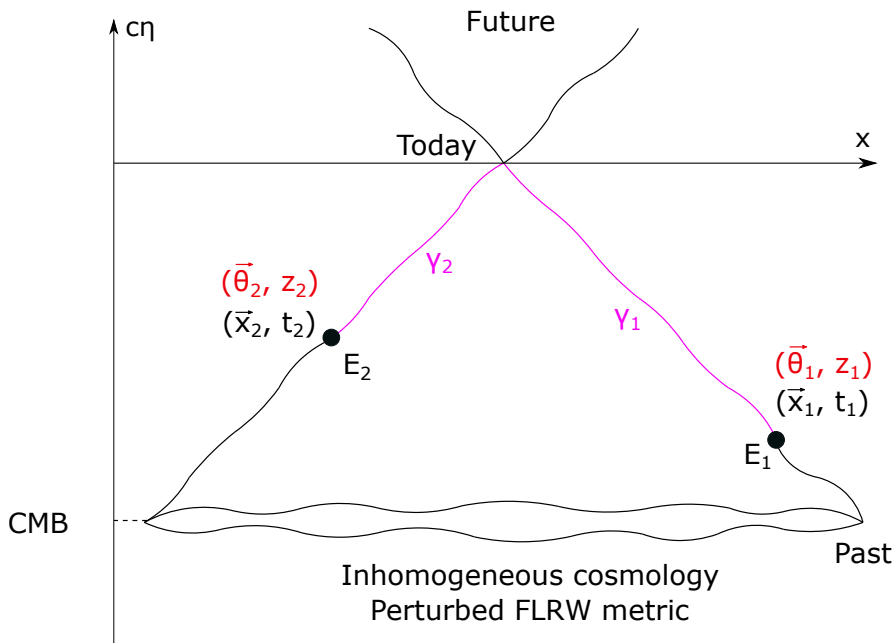


Photon propagation









These distortions blur the original signal BUT they also give informations about our universe!

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Problematic

How to properly map the '*real*' universe to the *observed* universe in order to extract relevant cosmological information?

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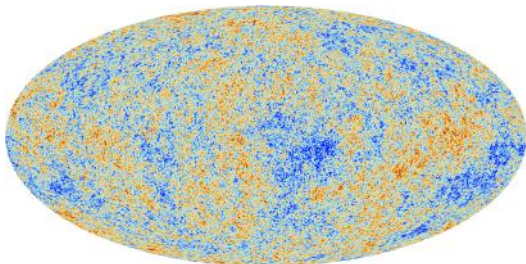
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Program

- Initial conditions → Matter dynamics → Structure formation (HPC Sim)
- Kinematics of photons → What we observe (HPC Sim)
- Statistical tools → Comparison to theory (Stats/Theory)
- Likelihood analysis & parameter constraints (Stats)

Numerical simulations



- Early universe
- Small density fluctuations
- Gaussian distribution

↘ Non-linear evolution ↘

- Late universe
- Virialised structures
- Non-gaussian distribution



N -body or Hydrodynamical simulation ?

N -body simulation

Compute the dynamics for dark matter particles

- Compute force and equation of motion
- Fast - Largest simulations with $\sim 16.364^3$ particles
- Physics well understood
- Not realistic at small scales (< 1 Mpc)

Hydro simulation

Compute the dynamics for dark matter and baryons

- Compute the moments of Boltzmann equation
- Slow - Largest simulations with $\sim 4096^3$ particles
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Different methods to compute the force

Particle-Particle (PP)

Direct pair-counting algorithm. Scales as $O(N^2)$

Particle-Mesh (PM)

Compute Poisson equation on a grid. Scales as $O(N \log N)$

Particle-Particle Particle-Mesh (P³M)

PP at short range, PM at large range. Scales as $O(N^2)$ and $N \log N$

Tree methods

Far enough particles (or cells) are aggregated. Scales as $O(N \log N)$

Fast-multipole method (FMM)

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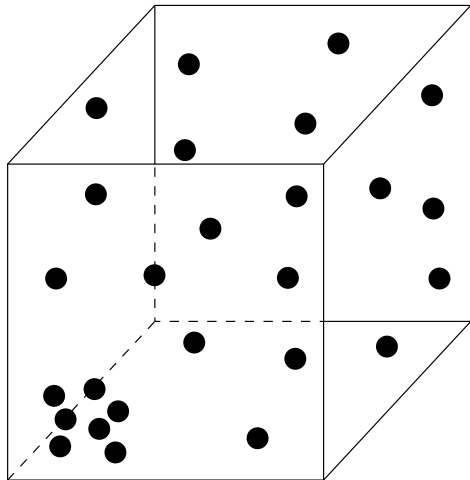
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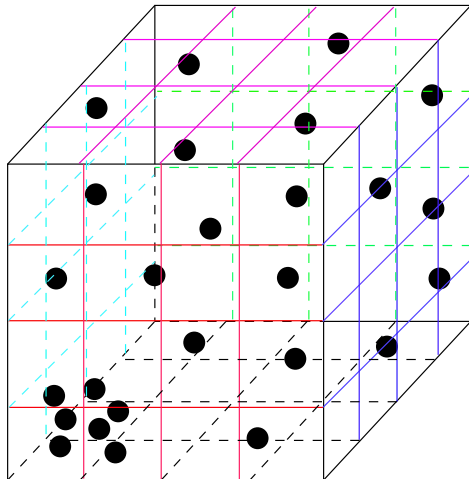
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Particle-Mesh (PM)



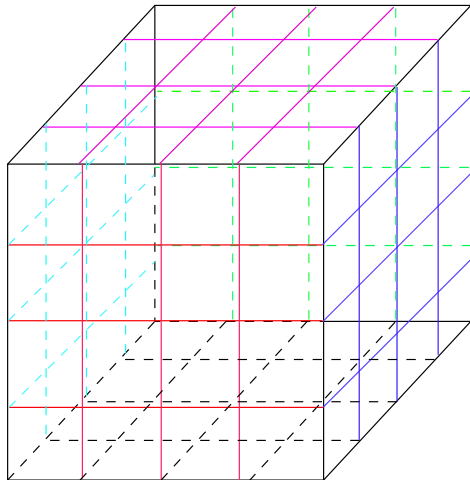
- Particle distribution in snapshot

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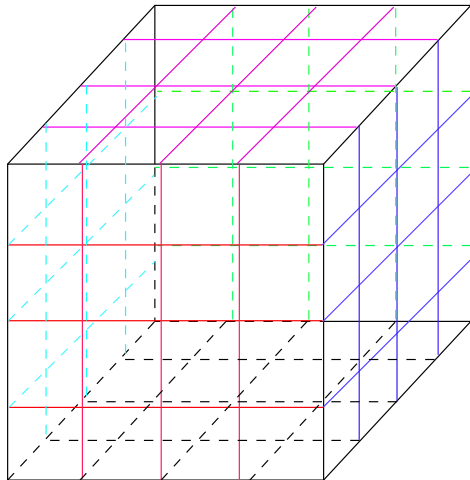
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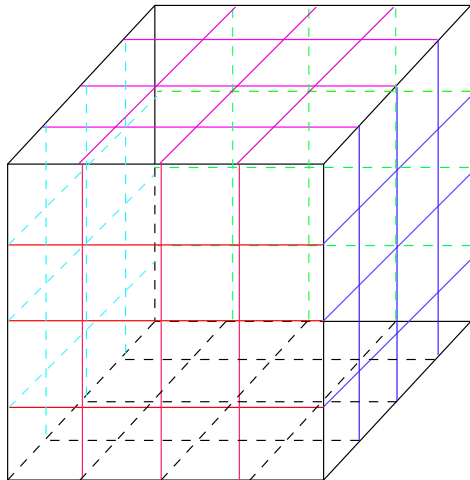
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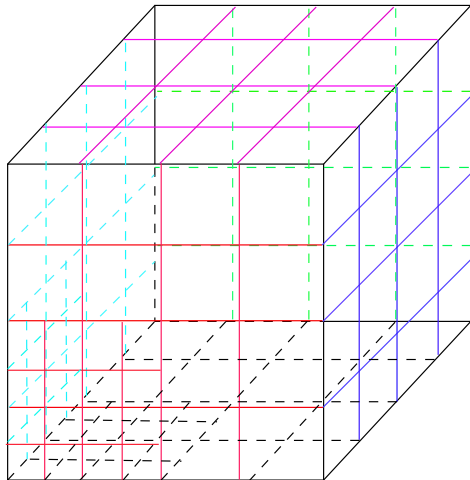
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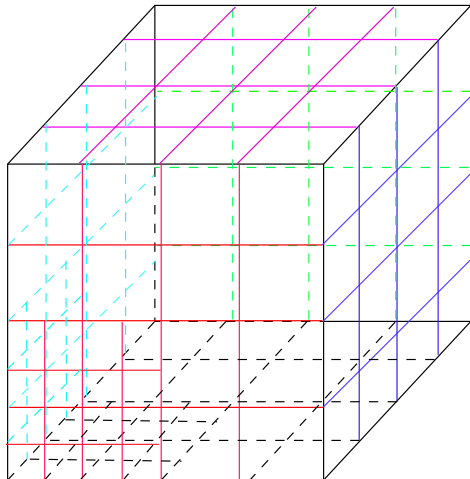
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- Equation of motion

Particle-Mesh with Adaptive Mesh Refinement (PM-AMR)



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- Compute density in cells
- $\Delta\Phi \propto \rho$
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- And same with AMR!
(multigrid method, Guillet & Teyssier 2011)

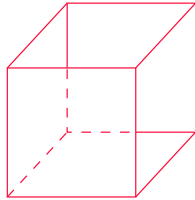
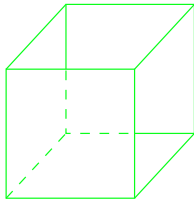
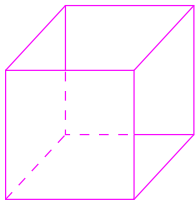
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(multigrid method, Guillet & Teyssier 2011)
- Particles : $\{\mathbf{x}_p, \mathbf{v}_p\}$
Cells : $\{\mathbf{x}_c, \Phi, \nabla_r\Phi, \rho\}$

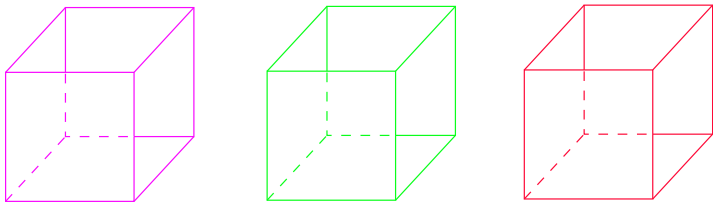
N -body simulation outputs

Snapshots : Particles + cells at some given constant redshifts (or time)

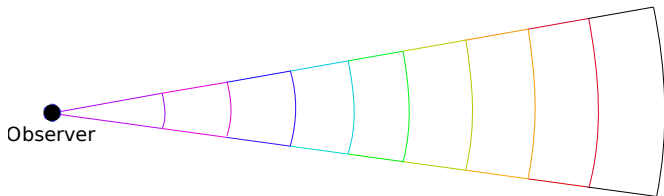


N -body simulation outputs

Snapshots : Particles + cells at some given constant redshifts (or time)



Lightcone : Particles + cells in concatenated shells from each step



Ray-tracing

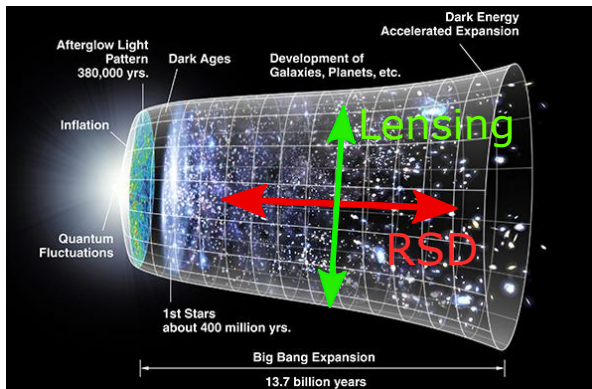
The inhomogeneous universe

Light propagation in a **perturbed** FLRW metric :

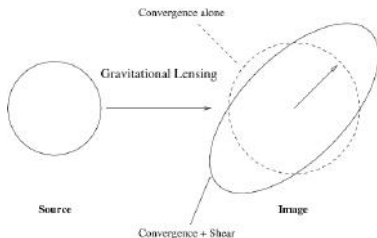
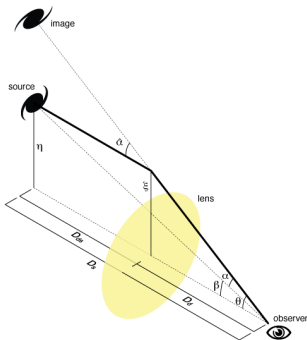
$$ds^2 = -(1+2\Psi)dt^2 + a(t)^2(1-2\Phi)[dx^2 + dy^2 + dz^2]$$

$1+z = \frac{(g_{\mu\nu}k^\mu k^\nu)_s}{(g_{\mu\nu}k^\mu k^\nu)_o} \neq$ homogeneous FLRW redshift (**Redshift-space distortions**)

apparent angles \neq true angles (**Lensing**)



Weak Lensing (WL)



Lens equation

$$\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

- $\vec{\beta}$ true angle
- $\vec{\theta}$ angle seen
- $\vec{\alpha}$ deflection angle

Approximation : $\vec{\alpha} \ll 1$

Distortion matrix

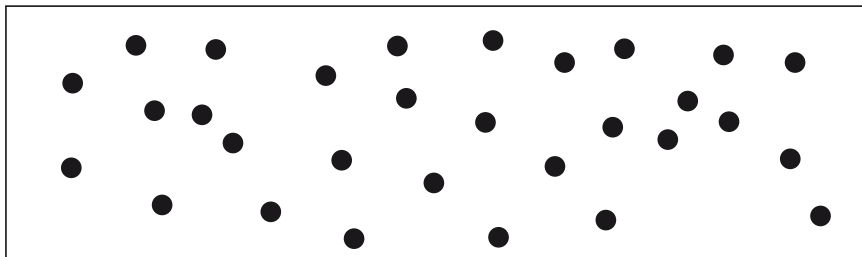
$$A_{ij} = \frac{\partial \beta_i}{\partial \theta_j} = \text{convergence} + \text{shear}$$

$$\text{convergence} : (1 - \kappa) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\text{shear} : -\gamma \begin{pmatrix} \cos 2\phi & \sin 2\phi \\ \sin 2\phi & -\cos 2\phi \end{pmatrix}$$

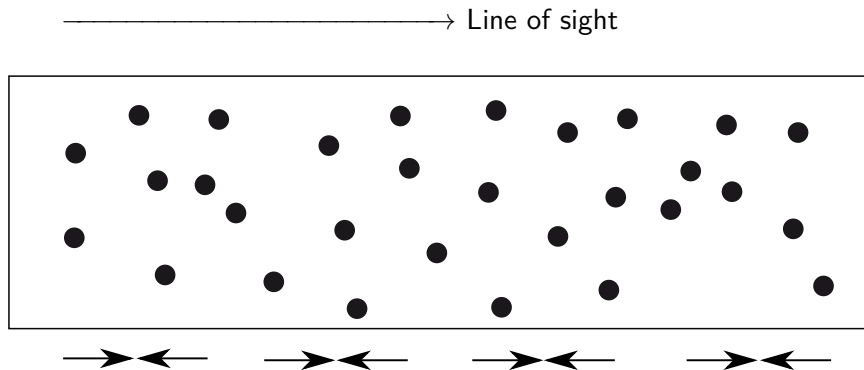
Standard approach for raytracing : The Born approximation

→ Line of sight



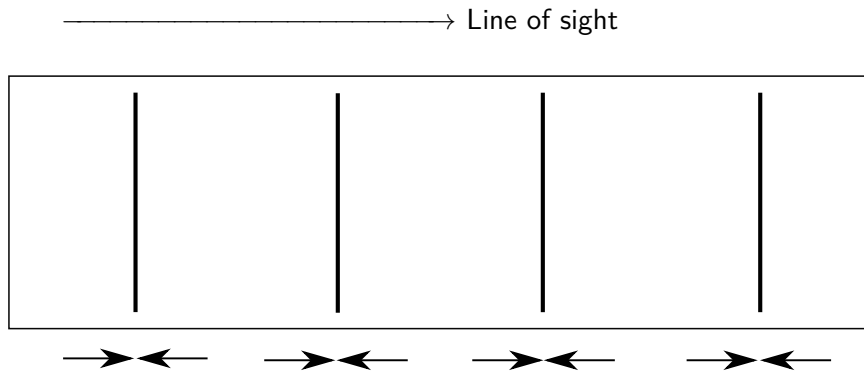
- Get the particle distribution in a lightcone pencil beam

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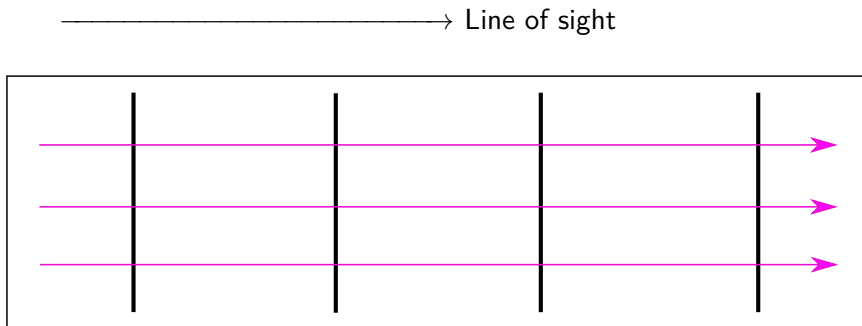
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- Project the density onto a plane (the lens!)
- Compute the Poisson equation to estimate the potential at the lens
- Sum the contributions of lenses along a given direction

Improved model : Post-Born

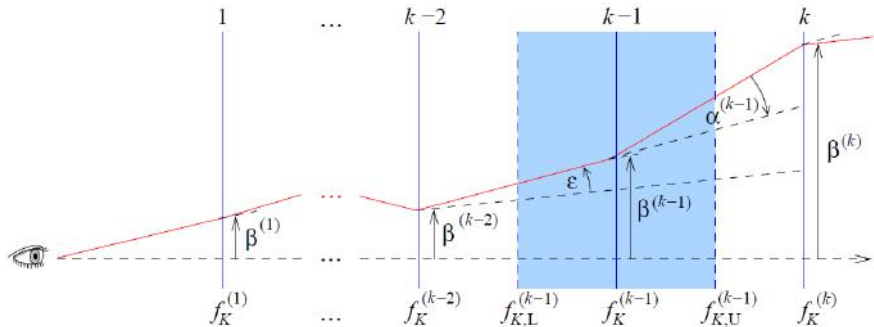


FIGURE – Hilbert et al. (2009)

Some limitations of standard approaches

- Approximations (Born, Post-Born, Lens equation, Flat sky...)
- Need to assume Ω_m to compute distances between lenses
- Restricted to “usual” weak lensing analysis

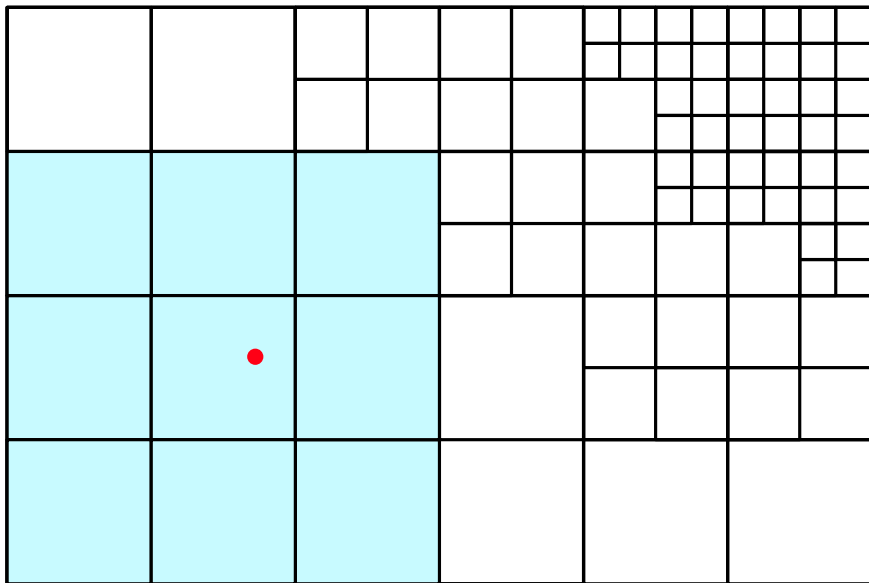
Ray-tracing : the MAGRATHEA-PATHFINDER code

A code for 3D AMR ray-tracing [Reverdy 2014 (thesis), MAB & Reverdy (2021)]

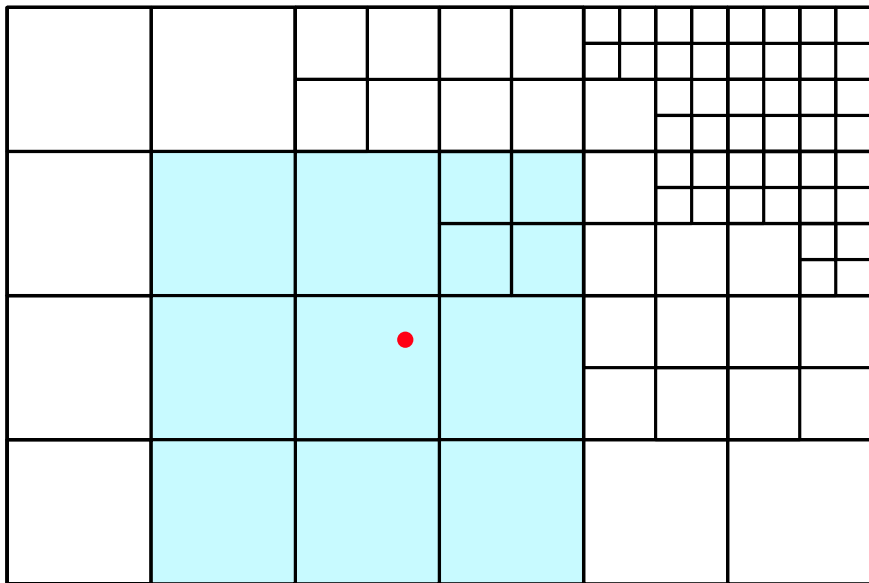
Our approach : [Direct geodesic integration of photons](#)

- C++11 Template Metaprogramming, MPI, multi-threading
- Backward integration using the geodesic equations
- RK4 integrator with 4 steps per AMR cell
- 3D interpolation with inverse TSC scheme

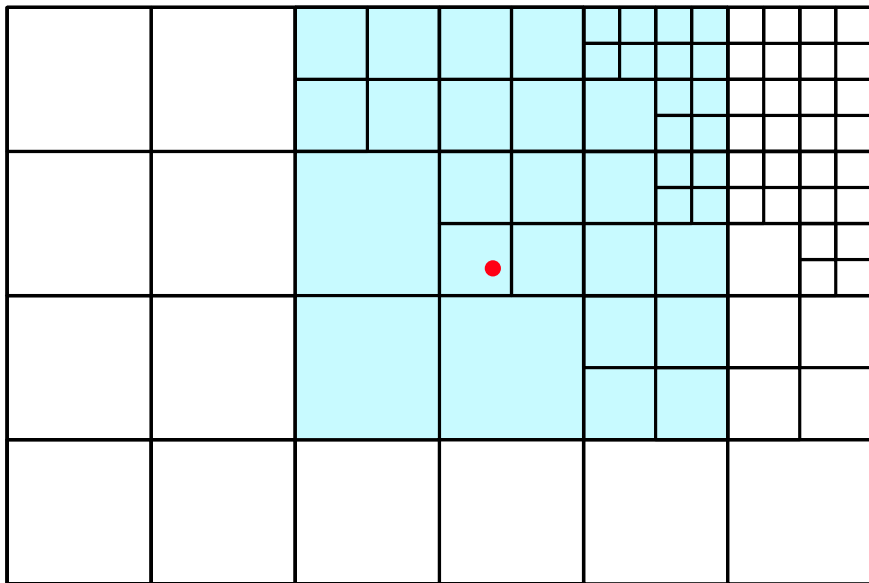
Integration and interpolation : TSC scheme [MAB & Reverdy (2021)]



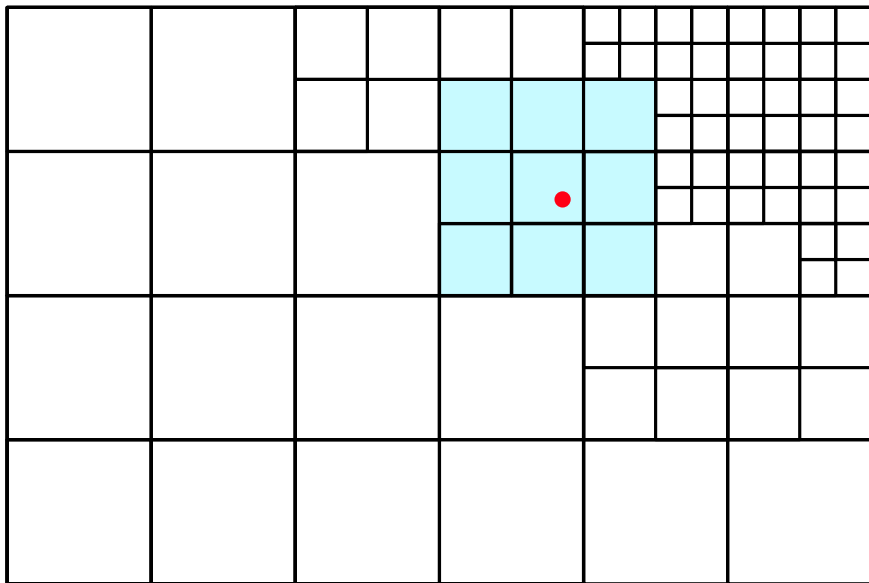
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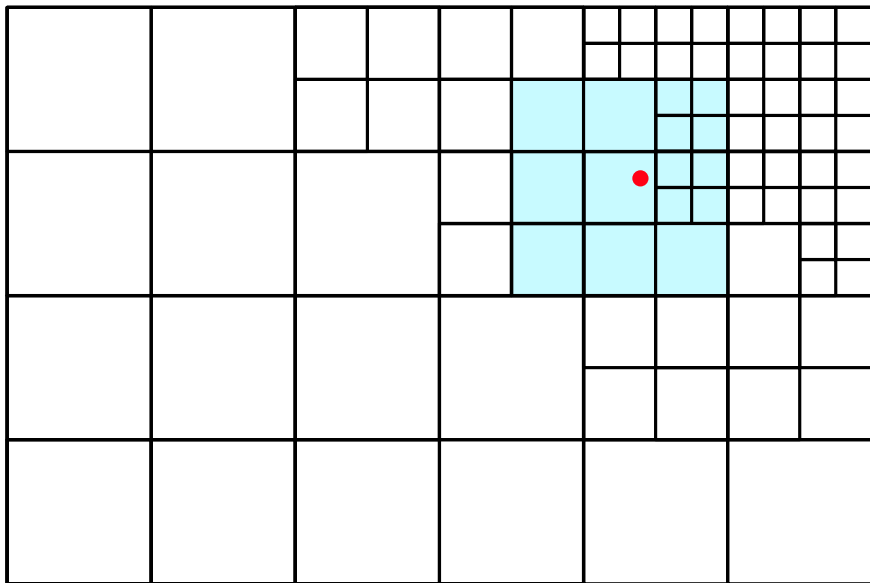
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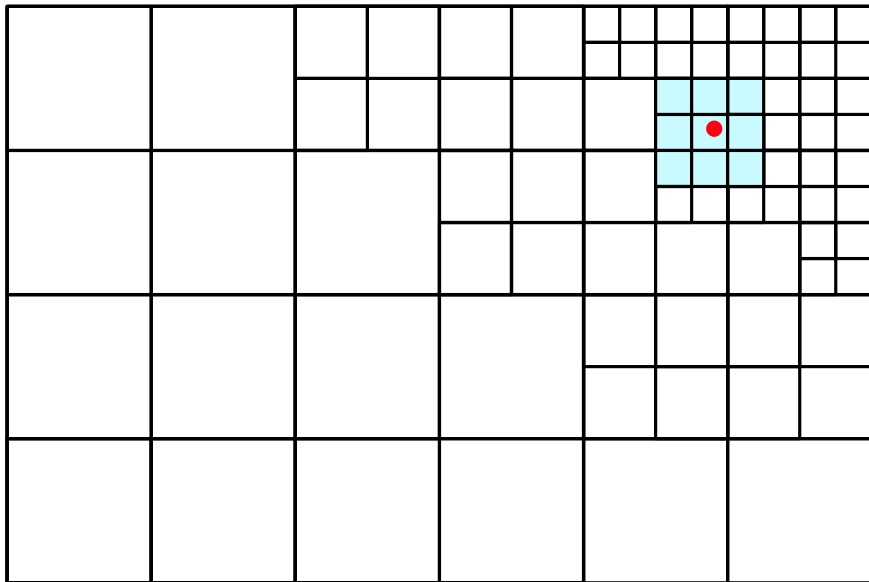
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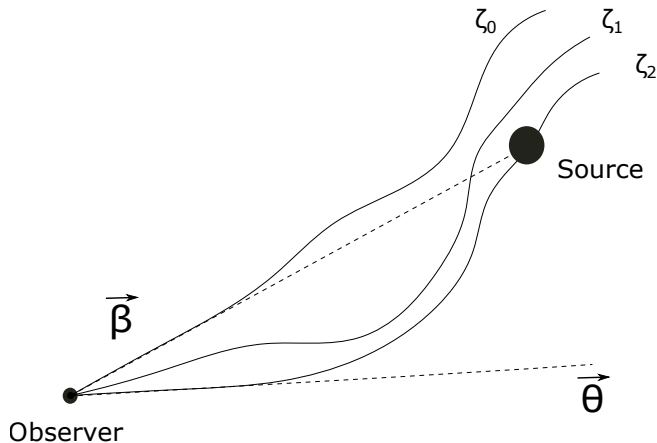


Geodesics finder

Null-geodesic finder

Find the connection between the observer and a source

→ Production of realistic observables! [MAB et al. (2019)]



Applications

Weak lensing with finite beams

Usually weak lensing assumes infinitesimal light beams. In reality, these beams are *finite*. Here we study the impact of beam finiteness on the convergence/shear.

The clustering of galaxies

Study the apparent distribution of sources beyond usual distortions due to peculiar velocities.

Simulation

The RayGal simulation suite [MAB et al. (2019), Rasera et al. (2021)]

Newtonian N-body simulations of interacting dark matter particles

Code : **RAMSES** [Teyssier 2002]

PM - AMR method

Gravity lightcone

RAYGALGROUPSIMS :

4096^3 particles, 2.625 Gpc/h box size

$M_{\text{part}} = 1.9 \times 10^{10} M_{\odot}$

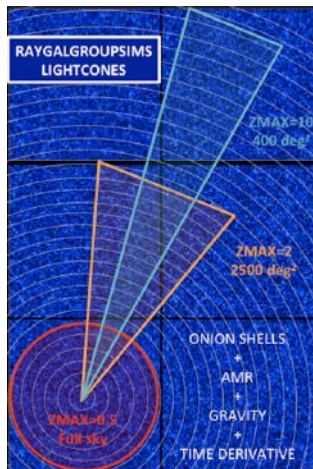
Lightcone halo detection : pFoF [Roy et al, 2014]

Initial conditions : **MPGRAFIC** [Prunet+08]

Calibrated on WMAP7 [Komatsu et al. 2011]

2 cosmologies :

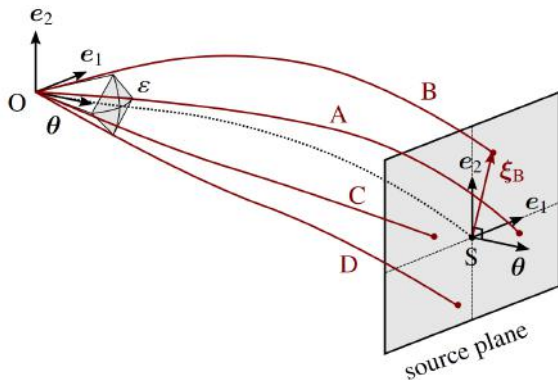
Λ CDM – w CDM



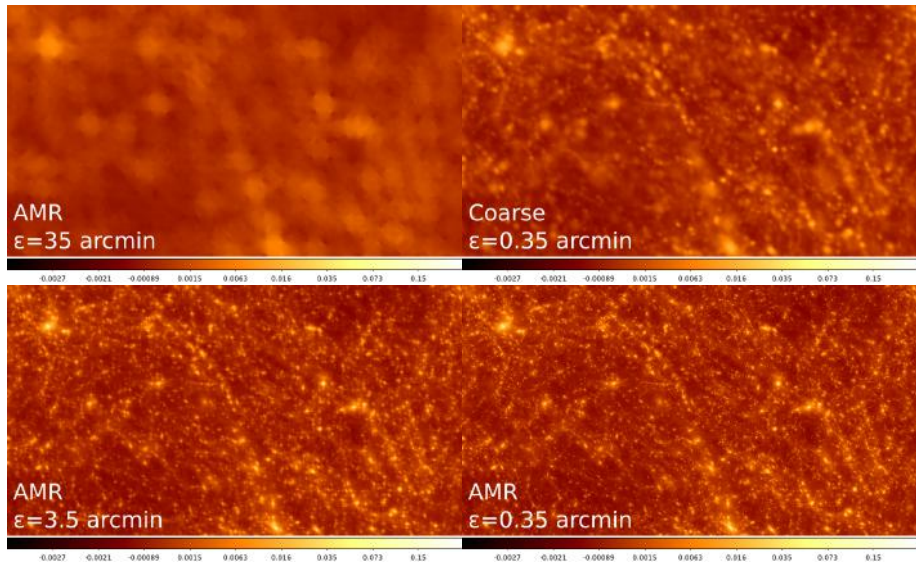
Weak lensing with finite beams

Geometry of the light beam [MAB & Fleury (2021)]

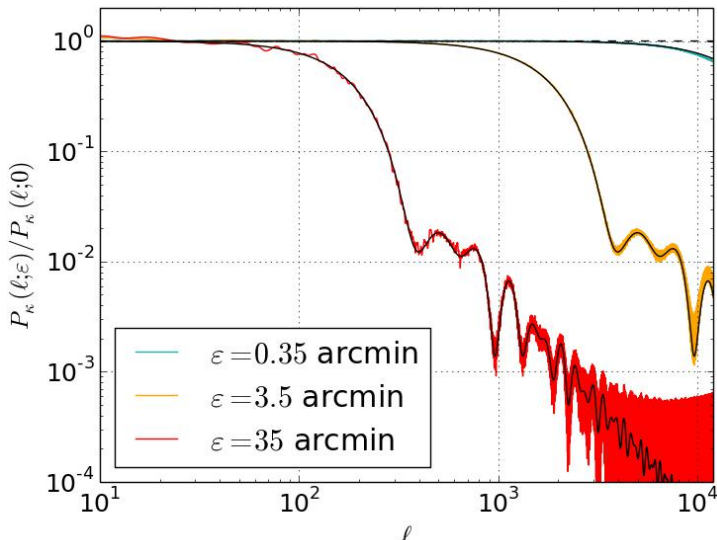
- Ray-bundle method
- Every photon follows $ds^2 = 0$
- Relevant for extended sources
- Compute $A_{ij} = \partial\beta_i/\partial\theta_j$



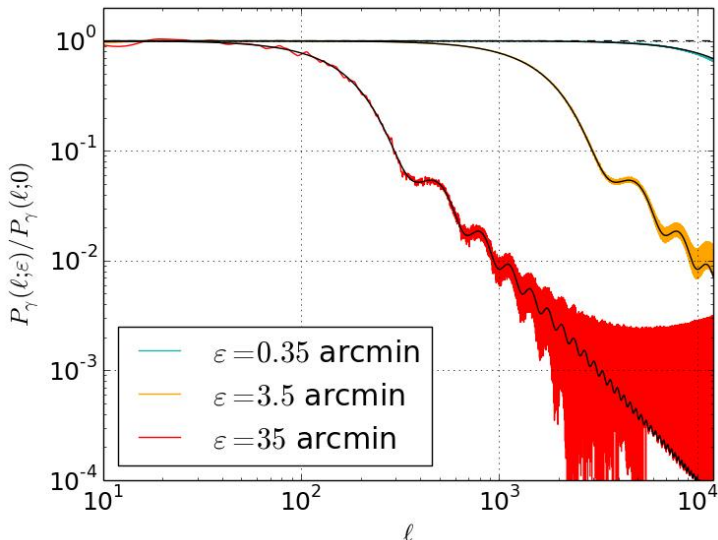
Convergence maps with finite beam [MAB & Fleury (2021)]



Convergence power spectrum with finite beam [MAB & Fleury (2021)]

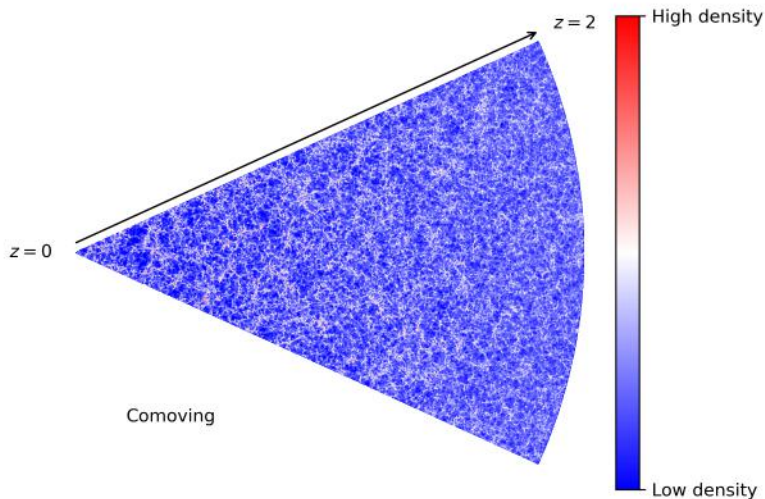


Shear power spectrum with finite beam [MAB & Fleury (2021)]

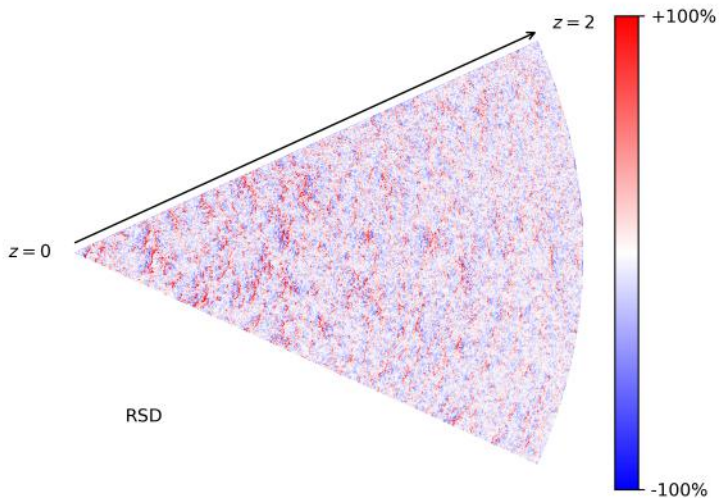


Galaxy clustering analysis

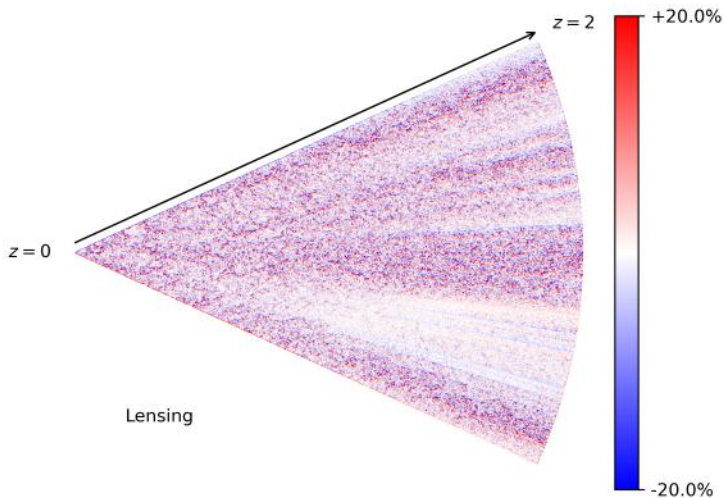
RayGal simulation light-cone [Rasera et al. (2021)]



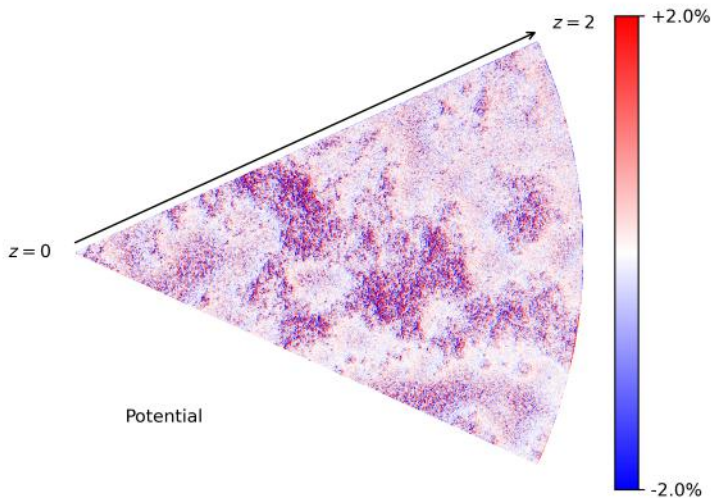
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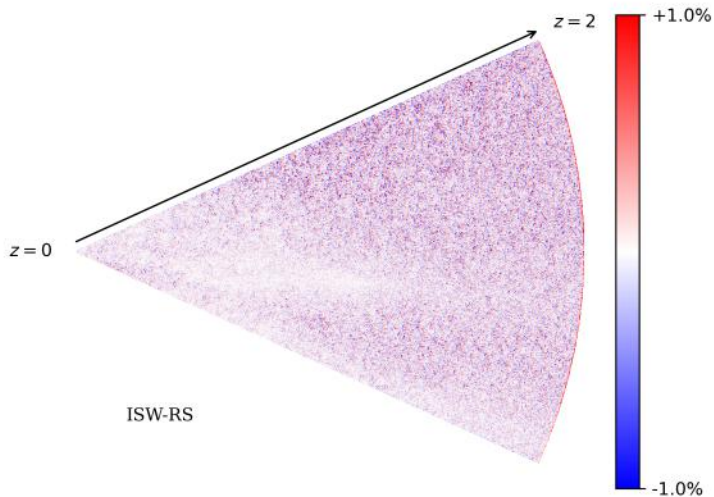
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Multipoles of the 2PCF in the RayGal simulation

redshift bins : **[0.8 – 1.0]** and **[1.6 – 1.9]**

ξ_ℓ	Doppler	v_o	Grav. redshift	Lensing*	T. Doppler	ISW
ξ_0	> 20%	3%	< 1%	1 – 10%	< 1%	< 1%
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* Angular displacement only

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Number counts heavily impacted by peculiar velocities (Doppler effect) and gravitational lensing through the change in observed angular positions and magnified fluxes!

Galaxy Clustering Redshift-Space Distortions

- Observed galaxy distribution modified by peculiar velocities
- Peculiar velocities depend on the growth rate of structures $f(z) = \Omega_m(z)^\gamma$
- Ω_m the amount of matter, $\gamma \sim 0.55$ in GR

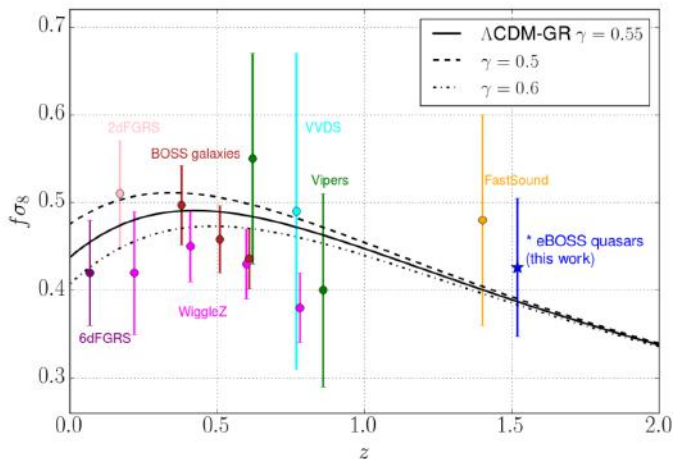
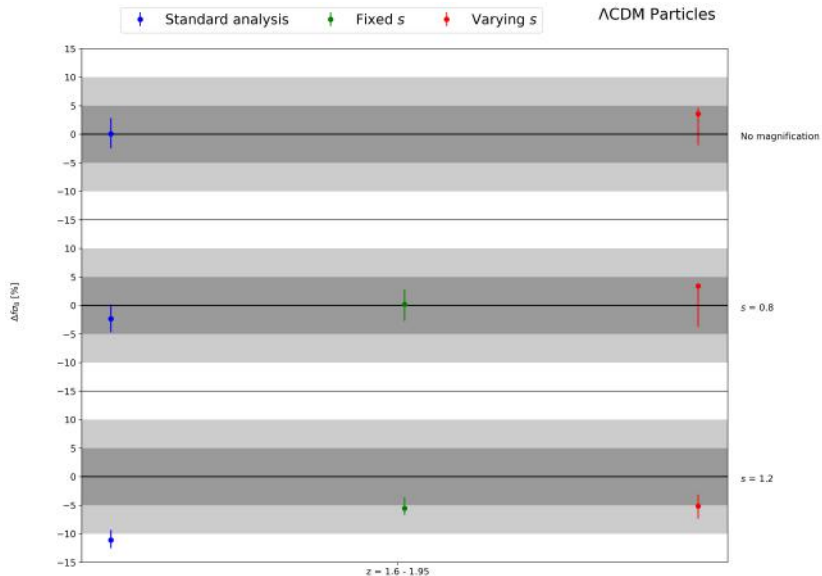


Figure from
Zarrouk et al. (2018)

Impact of lensing on the growth rate [MAB et al. (2021)]



Summary

An end-to-end pipeline

- Compute dynamics for the matter field using N-body simulations
- Extract lightcone from the simulation
- Perform ray-tracing on the lightcone to produce realistic observables
- Compare these observables to theoretical predictions
- Infer possible biases on cosmological parameter constraints
- Wide range of applications for our ray-tracing methods

Raytracing integration

$$ds^2 = a(t)^2 [-(1 + 2\Phi/c^2)d\eta^2 + (1 - 2\Phi/c^2)\delta_{ij}dx^i dx^j]$$

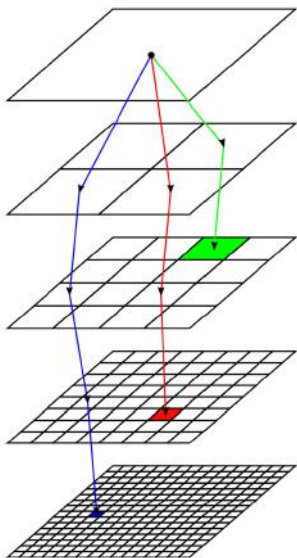
Geodesic equation

$$\frac{d^2\eta}{d\lambda^2} = -\frac{2a'}{a} \frac{d\eta}{d\lambda} \frac{d\eta}{d\lambda} - \frac{2}{c^2} \frac{d\phi}{d\lambda} \frac{d\eta}{d\lambda} + 2 \frac{\partial\phi}{\partial\eta} \left(\frac{d\eta}{d\lambda}\right)^2$$

$$\frac{d^2x}{d\lambda^2} = -\frac{2a'}{a} \frac{d\eta}{d\lambda} \frac{dx}{d\lambda} + \frac{2}{c^2} \frac{d\phi}{d\lambda} \frac{dx}{d\lambda} - 2 \frac{\partial\phi}{\partial x} \left(\frac{d\eta}{d\lambda}\right)^2$$

$$\frac{d^2y}{d\lambda^2} = -\frac{2a'}{a} \frac{d\eta}{d\lambda} \frac{dy}{d\lambda} + \frac{2}{c^2} \frac{d\phi}{d\lambda} \frac{dy}{d\lambda} - 2 \frac{\partial\phi}{\partial y} \left(\frac{d\eta}{d\lambda}\right)^2$$

$$\frac{d^2z}{d\lambda^2} = -\frac{2a'}{a} \frac{d\eta}{d\lambda} \frac{dz}{d\lambda} + \frac{2}{c^2} \frac{d\phi}{d\lambda} \frac{dz}{d\lambda} - 2 \frac{\partial\phi}{\partial z} \left(\frac{d\eta}{d\lambda}\right)^2$$



PROGRAM	CODE	INTEGER (BASE 2)	DECIMAL
<u>GOTO</u> UPPER RIGHT <u>GOTO</u> UPPER LEFT <u>STOP</u>	$\frac{1}{1}$ 11 $\frac{1}{1}$ 10 <u>0</u>	1111100000000000	63488
<u>GOTO</u> LOWER RIGHT <u>GOTO</u> UPPER LEFT <u>GOTO</u> LOWER RIGHT <u>STOP</u>	$\frac{1}{1}$ 01 $\frac{1}{1}$ 10 $\frac{1}{1}$ 01 <u>0</u>	1011101010000000	47744
<u>GOTO</u> LOWER LEFT <u>GOTO</u> UPPER LEFT <u>GOTO</u> UPPER RIGHT <u>GOTO</u> UPPER RIGHT <u>STOP</u>	$\frac{1}{1}$ 00 $\frac{1}{1}$ 10 $\frac{1}{1}$ 11 $\frac{1}{1}$ 11 <u>0</u>	1001101111110000	39920

FIGURE – MAGRATHEA indexing, from Reverdy (2014)

Redshift-space number count (linear) decomposition

$$\Delta^{\text{std}} = b\delta - \frac{1}{\mathcal{H}} \nabla_r(\mathbf{v} \cdot \mathbf{n}), \quad (1)$$

$$\Delta^{\text{acc}} = \frac{1}{\mathcal{H}c} \dot{\mathbf{v}} \cdot \mathbf{n}, \quad (2)$$

$$\Delta^{\text{q}} = -\frac{\dot{\mathcal{H}}}{c\mathcal{H}^2} \mathbf{v} \cdot \mathbf{n}, \quad (3)$$

$$\Delta^{\text{div}} = -\frac{2}{\mathcal{H}\chi} \mathbf{v} \cdot \mathbf{n}, \quad (4)$$

$$\Delta^{\text{pot,(1)}} = \frac{1}{\mathcal{H}c} \nabla_r \psi \cdot \mathbf{n}, \quad (5)$$

$$\Delta^{\text{pot,(2)}} = \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2c}{\mathcal{H}\chi} \right) \psi/c^2 - \frac{1}{\mathcal{H}c^2} \dot{\psi}, \quad (6)$$

$$\Delta^{\text{shapiro}} = (\phi + \psi)/c^2, \quad (7)$$

$$\Delta^{\text{lens}} = -\frac{1}{c^2} \int_0^{\chi} \frac{(\chi - \chi')\chi'}{\chi} \nabla_{\perp}^2 (\phi + \psi) d\chi', \quad (8)$$

$$\Delta^{\text{isw}} = \frac{1}{\mathcal{H}c^2} (\dot{\phi} + \dot{\psi}), \quad (9)$$

$$\Delta^{\text{LC}} = \mathbf{v} \cdot \mathbf{n}/c, \quad (10)$$

$$\Delta^{\text{neglect}} = \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2c}{\mathcal{H}\chi} \right) \frac{1}{c^2} \int_{\eta}^{\eta_0} \frac{\partial(\phi + \psi)}{\partial\eta} d\eta' + \frac{2}{\chi c^2} \int_0^{\chi} (\phi + \psi) d\chi'. \quad (11)$$